# Linear reducibility of Feynman integrals

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Some FI are expressible via multiple polylogarithms (MPL)

$$\operatorname{Li}_{n_1,\ldots,n_d}(z_1,\ldots,z_d) = \sum_{0 < k_1 < \cdots < k_d} \frac{z_1^{k_1} \cdots z_d^{k_d}}{k_1^{n_1} \cdots k_d^{n_d}}$$

Example 
$$(p_1^2 = 1, p_2^2 = z\bar{z}, p_3^2 = (1 - z)(1 - \bar{z}))$$
  
 $\Phi\left( \underbrace{}^{+} \underbrace{}^{+} \right) = \frac{4i \ln [\text{Li}_2(z) + \log(1 - z) \log |z|]}{z - \bar{z}}$ 

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## Example (massless propagators)

$$\Phi\left(-\underbrace{-}\right) = 6\zeta_{3}, \quad \Phi\left(-\underbrace{-}\right) = 252\zeta_{3}\zeta_{5} + \frac{432}{5}\zeta_{3,5} - \frac{25056}{875}\zeta_{2}^{4}$$

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#### Questions

- What comes beyond MPL? [Weinzierl's talk]
- I How to tell if a FI evaluates to MPL? What is the alphabet?
- O How to compute it explicitly?

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- I How to compute it explicitly in an automated way?

#### integral representations

- momentum space [IBP, unitarity]
- position space [Gardi's talk, Drummond & Schnetz graphical functions]
- twistor space [Trnka's talk]
- parametric space [this talk]
- Mellin-Barnes
- sum representations [Schneider's talk]
- differential equations [Smirnov's talk]

# Tools to study Feynman integrals

#### integral representations

- momentum space [IBP, unitarity]
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## Plan of the talk

- predicting bounds on the singularities of an integral
- nice (linearly reducible) integral representations can be integrated in terms of MPL
- good parametrizations for FI via Laplace transform

# Definite integral representations

• hypergeometric functions  ${}_{p}F_{q}(\cdots;z)$ 

$${}_{2}F_{1}\left(\begin{array}{c}a,b\\c\end{array}\right|z\right) = \frac{\Gamma(c)}{\Gamma(b)\Gamma(c-b)}\int_{0}^{1}t^{b-1}(1-t)^{c-b-1}(1-zt)^{-a}\mathrm{d}t$$

• Appell's functions  $F_1$ ,  $F_2$ ,  $F_3$ ,  $F_4$ 

$$F_{3}\left(\begin{array}{c}\alpha,\alpha'\\\beta,\beta'\end{array}\middle|\gamma\middle|x\ y\right) = \frac{\Gamma(\gamma)}{\Gamma(\beta)\Gamma(\beta')\Gamma(\gamma-\beta-\beta')}$$
$$\times \int_{0}^{1}\int_{0}^{1-\nu}u^{\beta-1}v^{\beta'-1}(1-u-\nu)^{\gamma-\beta-\beta'-1}(1-ux)^{-\alpha}(1-\nu y)^{-\alpha'}\mathrm{d}u\mathrm{d}v$$

- Feynman integrals in Schwinger parameters
- Phase-space integrals

• . . .

With the superficial degree of divergence  $sdd = |E(G)| - D/2 \cdot loops(G)$ ,

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Graph polynomials:

$$\psi = \sum_{T} \prod_{e \notin T} \alpha_{e} \qquad \varphi = \sum_{F = T_{1} \cup T_{2}} q^{2} (T_{1}) \prod_{e \notin F} \alpha_{e} + \psi \sum_{e} m_{e}^{2} \alpha_{e}$$

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Example (D = 4)  $p_1 \xrightarrow{3} p_2 \qquad \psi =$   $G = 2 \xrightarrow{1} \qquad \varphi =$  $p_3 \qquad \Phi(G) = \iint \frac{\mathrm{d}\alpha_2 \,\mathrm{d}\alpha_3}{\psi \varphi}\Big|_{\alpha_1 = 1}$ 

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$$G = \begin{array}{c} p_1 & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & &$$

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$$G = \begin{array}{c} p_1 & & 3 \\ \varphi = \\ p_3 \end{array} \qquad \qquad \psi = \alpha_1 + \alpha_2 + \alpha_3 \\ \varphi = \\ \varphi = \\ \varphi = \\ \varphi(G) = \iint \left. \frac{\mathrm{d}\alpha_2 \, \mathrm{d}\alpha_3}{\psi \varphi} \right|_{\alpha_1 = 1}$$

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Singularities of the original integrand:  $S = \{\psi, \varphi\}$ , i.e. at  $\alpha_3 = \sigma_i$  for

$$\sigma_1 = -1 - \alpha_2$$
 and  $\sigma_2 = -\frac{lpha_2(1-z)(1-ar z)}{lpha_2 + zar z}$ 

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After integrating  $\alpha_1$  from 0 to  $\infty,$  the integrand has singularities

$$S_{3} = \{\underbrace{1+\alpha_{2}}_{\sigma_{1}=0}, \underbrace{\alpha_{2}, 1-z, 1-\bar{z}}_{\sigma_{2}=0}, \underbrace{\alpha_{2}+z\bar{z}}_{\sigma_{2}=\infty}, \underbrace{z+\alpha_{2}, \bar{z}+\alpha_{2}}_{\sigma_{1}=\sigma_{2}}, \underbrace{z+\alpha_{2}, \bar{z}+\alpha_{2}}_{\sigma_{2}=\sigma_{2}}, \underbrace{z+\alpha_{2}, \bar{z}+\alpha_{2}}, \underbrace{z+\alpha_$$

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$$S_{3} = \left\{ \underbrace{1 + \alpha_{2}}_{\sigma_{1} = 0}, \underbrace{\alpha_{2}, 1 - z, 1 - \bar{z}}_{\sigma_{2} = 0}, \underbrace{\alpha_{2} + z\bar{z}}_{\sigma_{2} = \infty}, \underbrace{z + \alpha_{2}, \bar{z} + \alpha_{2}}_{\sigma_{1} = \sigma_{2}} \right\}$$

With the same logic, predict the possible singularities after  $\int_0^\infty d\alpha_2$ :

$$S_{3,2} = \{z, \bar{z}, 1-z, 1-\bar{z}, z-\bar{z}, z\bar{z}-1\}$$

$$\Phi\left(\overbrace{-}^{\bullet}\right) = \int \frac{\mathrm{d}\alpha_2 \,\mathrm{d}\alpha_3}{(1+\alpha_2+\alpha_3)(\alpha_2\alpha_3+z\bar{z}\alpha_3+(1-z)(1-\bar{z})\alpha_2)}$$
$$= \frac{1}{z-\bar{z}} \int \left(\frac{\mathrm{d}\alpha_2}{\alpha_2+\bar{z}}-\frac{\mathrm{d}\alpha_2}{\alpha_2+z}\right) \log \frac{(\alpha_2+1)(\alpha_2+z\bar{z})}{(1-z)(1-\bar{z})\alpha_2}$$

Singularities of the original integrand:  $S = \{\psi, \varphi\}$ , i.e. at  $\alpha_3 = \sigma_i$  for

$$\sigma_1 = -1 - \alpha_2$$
 and  $\sigma_2 = -\frac{\alpha_2(1-z)(1-\overline{z})}{\alpha_2 + z\overline{z}}$ 

After integrating  $\alpha_1$  from 0 to  $\infty$ , the integrand has singularities

$$S_{3} = \left\{ \underbrace{1 + \alpha_{2}}_{\sigma_{1} = 0}, \underbrace{\alpha_{2}, 1 - z, 1 - \bar{z}}_{\sigma_{2} = 0}, \underbrace{\alpha_{2} + z\bar{z}}_{\sigma_{2} = \infty}, \underbrace{z + \alpha_{2}, \bar{z} + \alpha_{2}}_{\sigma_{1} = \sigma_{2}} \right\}$$

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# Polynomial reduction [F. Brown]

## Definition

Let S denote a set of polynomials, then  $S_e$  are the irreducible factors of

$$\left\{ \mathsf{lead}_e(f), \left. f \right|_{lpha_e = 0}, D_e(f) \colon \ f \in S 
ight\} \quad \mathsf{and} \quad \left\{ [f,g]_e \colon \ f,g \in S 
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If the singularities of F are cointained in S, then the singularities of  $\int_0^\infty F d\alpha_e$  are contained in S<sub>e</sub>.

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This gives only very coarse upper bounds. For example,  $z\overline{z} - 1$  is spurious: It drops out in  $S_{2,3} \cap S_{3,2} = \{z, \overline{z}, 1 - z, 1 - \overline{z}, z - \overline{z}\}$  because

$$S_{2,3} = \{z, \bar{z}, 1-z, 1-\bar{z}, z-\bar{z}, z\bar{z}-z-\bar{z}\}.$$

#### Improvements

- Fubini: intersect over different orders
- Compatibility graphs

# Compatibility graphs

Keep track of compatibilities  $C \subset {S \choose 2}$  between polynomials

- $\bullet\,$  for  $S_e,$  we only need the resultants  $[f,g]_e$  for compatible  $\{f,g\}\in C$
- only pairs  $[f,g]_e$ ,  $[g,h]_e$  (with one polynomial in common) become compatible in  $C_e$



Since  $z\bar{z}\alpha_1 + \alpha_2$  and  $\alpha_1 + \alpha_2$  are not compatible, their resultant  $1 - z\bar{z}$  does not occur.

# HyperInt: massive triangle

Graph polynomials:

- > E:=[[2,3],[1,3],[1,2]]:
- > M:=[[1,s1],[2,s2],[3,s3]],[m1<sup>2</sup>,m2<sup>2</sup>,m3<sup>3</sup>]:
- > psi:=graphPolynomial(E):
- > phi:=secondPolynomial(E,M):

Polynomial reduction:

1

- > L[{}]:=[{psi,phi}, {{psi,phi}}]:
- > cgReduction(L, {s1,s2,s3,m1,m2,m3}, 2):
- > L[{x[1],x[2]}][1];

$$\left\{s_i + (m_j \pm m_k)^2, \sum_i s_i^2 - \sum_{i \neq j} s_i s_j\right\}$$

$$\cup \left\{ s_1 s_2 s_3 - \sum_i s_i^2 m_i^2 + \sum_i s_i (m_i^2 - m_j^2) (m_i^2 - m_k^2) + \sum_{i < j} s_i s_j (m_i^2 + m_j^2) \right\}$$

#### Definition

If for some order of variables (edges), all  $S_{1,...,k}$  are linear in  $\alpha_{k+1}$ , then S (the Feynman graph G with  $S = \{\psi, \varphi\}$ ) is called linearly reducible.

#### Lemma

If S is linearly reducible, the integral  $\prod_e \int_0^\infty d\alpha_e f$  of any rational function f with singularities in S is a MPL with symbol letters in  $S_{1,...,N}$ .

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Integration of linearly reducible integrands in terms of MPL can be automatized via hyperlogarithms.

Definition (Hyperlogarithms, Lappo-Danilevsky 1927)

$$G(\sigma_1,\ldots,\sigma_w;z) := \int_0^z \frac{\mathrm{d} z_1}{z_1 - \sigma_1} G(\sigma_2,\ldots,\sigma_w;z_1)$$

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Problem: Given  $G(\vec{\sigma}(\alpha); z)$ , write it as hyperlogarithm with constant letters and final argument  $\alpha$ . Recursive solution via

$$dG(\vec{\sigma};z) = \sum_{i=1}^{n} G(\cdots, \phi_i, \cdots; z) d\log \frac{\sigma_i - \sigma_{i-1}}{\sigma_i - \sigma_{i+1}} \qquad \sigma_0 := z, \sigma_{n+1} := 0$$

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$$\frac{\partial}{\partial \alpha} G(0, -\alpha; 1) = -\frac{1}{\alpha} [G(0; \alpha) - G(-1; \alpha)]$$
  

$$\Rightarrow G(0, -\alpha; 1) = -G(0, 0; \alpha) + G(0, -1; \alpha)$$

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#### Example

$$\begin{aligned} &\frac{\partial}{\partial \alpha} G\left(0, -\alpha; 1\right) = -\frac{1}{\alpha} \left[ G(0; \alpha) - G(-1; \alpha) \right] \\ &\Rightarrow G\left(0, -\alpha; 1\right) = -G(0, 0; \alpha) + G(0, -1; \alpha) + \zeta_2 \end{aligned}$$

Computation of integration constants relies on shuffle algebra

$$G(\vec{\sigma};z) \cdot G(\vec{\tau};z) = G(\vec{\sigma} \sqcup \vec{\tau};z).$$

All this is implemented in the Maple program HyperInt. Similar things can be done with other programs.

# HyperInt

- open source Maple program
- integration of hyperlogarithms
- transformations of MPL to G(···; z)
   Note: No further simplification (e.g. rewrite as Li<sub>2,2</sub>) provided!
- polynomial reduction
- graph polynomials
- symbolic computation of constants (no numerics)

## Example

- > read "HyperInt.mpl":
- > hyperInt(polylog(2,-x)\*polylog(3,-1/x)/x,x=0..infinity):

 $\frac{8}{7}\zeta_2^3$ 

> fibrationBasis(%);

computes  $\int_0^\infty \operatorname{Li}_2(-x) \operatorname{Li}_3(-1/x) dx = \frac{8}{7}\zeta_2^3$ .

## HyperInt: propagator



- > E := [[2,1],[2,3],[2,5],[5,1],[5,3],[5,4],[4,1],[4,3]]: > psi := graphPolynomial(E):
- > phi := secondPolynomial(E, [[1,1], [3,1]]):
- > add((epsilon\*log(psi^5/phi^4))^n/n!,n=0..2)/psi^2:
- > hyperInt(eval(%,x[8]=1), [seq(x[n],n=1..7)]):
- > collect(fibrationBasis(%), epsilon);

$$\begin{split} \left( 254\zeta_7 + 780\zeta_5 - 200\zeta_2\zeta_5 - 196\zeta_3^2 + 80\zeta_2^3 - \frac{168}{5}\zeta_2^2\zeta_3 \right) \varepsilon^2 \\ + \left( -28\zeta_3^2 + 140\zeta_5 + \frac{80}{7}\zeta_2^3 \right) \varepsilon + 20\zeta_5 \end{split}$$
## HyperInt: triangle

Graph polynomials:

> E:=[[1,2],[2,3],[3,1]]:

- > M:=[[3,1],[1,z\*zz],[2,(1-z)\*(1-zz)]]:
- > psi:=graphPolynomial(E):
- > phi:=secondPolynomial(E,M):

Integration:

- > hyperInt(eval(1/psi/phi,x[3]=1),[x[1],x[2]]):
- > factor(fibrationBasis(%,[z,zz]));

 $\left(\mathsf{Hlog}\left(1;z\right)\mathsf{Hlog}\left(0;zz\right)-\mathsf{Hlog}\left(0;z\right)\mathsf{Hlog}\left(1;zz\right)+\mathsf{Hlog}\left(0,1;zz\right)\right.$ 

- Hlog(1, 0; zz) + Hlog(1, 0; z) - Hlog(0, 1; z))/(z - zz)

Polynomial reduction:

- > L[{}]:=[{psi,phi}, {{psi,phi}}]:
- > cgReduction(L):
- >  $L[{x[1],x[2]}][1];$

 $\{-1+z, -1+zz, -zz+z\}$ 

• all  $\leq$  4 loop massless propagators (Panzer)



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2 all  $\leq$  3 loop massless off-shell 3-point (Chavez & Duhr, Panzer)



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3 all  $\leq$  3 loop massless off-shell 3-point (Chavez & Duhr, Panzer)



**(**) all  $\leq$  2 loop massless on-shell 4-point (Lüders)



# Linearly reducible massive graphs (some examples)



• 3-constructible graphs (3-point functions) [Brown, Schnetz, Panzer]





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#### Theorem (Panzer)

All  $\epsilon$ -coefficients of these graphs (off-shell) are MPL over the alphabet  $\{z, \overline{z}, 1-z, 1-\overline{z}, z-\overline{z}, 1-z\overline{z}, 1-z-\overline{z}, z\overline{z}-z-\overline{z}\}.$ 

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• minors of ladder-boxes (up to 2 legs off-shell)



Theorem (Panzer)

All  $\epsilon$ -coefficients of these graphs are MPL. For the massless case, the alphabet is just  $\{x, 1 + x\}$  for x = s/t.

## 4-point recursions

Start with the box and repeat, in any order:

• Appending a vertex:



• Adding an edge:





Let  $f_3$ ,  $f_4$ ,  $f_{12}$  and  $f_{14}$  denote the spanning forest polynomials such that

$$\varphi = \mathcal{F} = (p_1 + p_2)^2 f_{12} + (p_1 + p_4)^2 f_{14} + p_3^2 f_3 + p_4^2 f_4$$

$$\begin{array}{c} v_{1} & v_{4} \\ 1 & 3 \\ v_{2} & v_{3} \end{array} \psi = \alpha_{1} + \alpha_{2} + \alpha_{3} + \alpha_{4} \quad f_{12} = \alpha_{2}\alpha_{4} \\ f_{14} = \alpha_{1}\alpha_{3} \end{array} \quad \begin{array}{c} f_{3} = \alpha_{2}\alpha_{3} \\ f_{4} = \alpha_{3}\alpha_{4} \end{array}$$

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#### Definition

$$F(G;z) := \int_{\mathbb{R}^{E}_{+}} \psi_{G}^{-D/2} \cdot \delta^{(4)} \left(\frac{f}{\psi} - z\right) \prod_{e \in E} \mathrm{d}\alpha_{e}^{a_{e}-1} \alpha_{e} \qquad (\mathbb{R}^{4}_{+} \longrightarrow \mathbb{R}_{+})$$

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$$F\begin{pmatrix}v_{1} & v_{4} \\ 1 & 3 \\ v_{2} & v_{3} \end{pmatrix} = \begin{cases} \frac{1}{z_{3}z_{4}} & (D=4) \\ \frac{z_{12}}{[z_{12}(z_{14}+z_{3}+z_{4})+z_{3}z_{4}]^{2}} \\ \frac{z_{12}}{Q} & (D=6) \end{cases}$$

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### Definition

$$\Phi(G) = \frac{\Gamma(\mathsf{sdd})}{\prod_e \Gamma(a_e)} \int_0^\infty \frac{F(G; z) \Omega}{\left[ (p_1 + p_2)^2 z_{12} + (p_1 + p_4)^2 z_{14} + p_3^2 z_3 + p_4^2 z_4 \right]^{\mathsf{sdd}}}$$

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$$G = \bigvee_{v_2} \bigvee_{v_3} \bigvee_{v_3} G' = \bigvee_{v_2} \bigvee_{v_2} \bigvee_{v_3} \bigvee_{v_2} \bigvee_{v_2} \bigvee_{v_3} \bigvee_{v_2} \bigvee_{v_3} \bigvee_{v_2} \bigvee_{v_3} \bigvee_{v_2} \bigvee_{v_3} \bigvee_{v_2} \bigvee_{v_3} \bigvee_{$$

## Example (D = 6 and $a_e = 1$ )

$$F\left(\bigcup_{v_{2}}^{v_{1}} \bigcup_{v_{3}}^{v_{4}}; z\right) = \int_{0}^{z_{3}} F\left(\bigcup_{v_{2}}^{v_{1}} \bigcup_{v_{3}}^{4}; z_{12}, z_{14}, z_{3}', z_{4}\right) dz_{3}'$$

11

$$G = \bigvee_{v_2}^{v_1} \bigvee_{v_3}^{v_4} \mapsto G' = \bigvee_{v_2}^{v_1} \bigvee_{v_3}^{v_4} \bigvee_{v_2}^{v_4} \bigvee_{v_2}^{v_4} \bigvee_{v_2}^{v_4} \bigvee_{v_2}^{v_4} \bigvee_{v_2}^{v_4} \bigvee_{v_3}^{v_4} \bigvee_{v_2}^{v_4} \bigvee_{v_3}^{v_4}$$
  
sing  $(f'_{12}, f'_{14}, f'_3, f'_4, \psi') = (f_{12}, f_{14}, f_3, f_4 + \alpha_e \psi, \psi),$   
 $F(G'; z) = \int_0^{z_4} F(G; z_{12}, z_{14}, z_3, z_4 - \alpha_e) \alpha_e^{a_e - 1} d\alpha_e$ 

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$$F\left(\bigcup_{v_{2}}^{v_{1}} \bigcup_{v_{3}}^{v_{4}}; z\right) = \int_{0}^{z_{3}} \frac{z_{12} dz'_{3}}{[z_{12}(z_{14} + z'_{3} + z_{4}) + z'_{3}z_{4}]^{2}} = \frac{z_{3}}{(z_{14} + z_{4}) \cdot Q}$$

11

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#### Example $(D = 6 \text{ and } a_e = 1)$



# Adding an edge



$$F_{G'}(z) = Q^{a_e + \mathsf{sdd} - D} \int_0^{z_{12}} x^{D/2 - 2} \left[ Q^{D/2 - \mathsf{sdd}} \cdot F_G \right]_{z_{12} = z_{12} - x} \mathrm{d}x$$

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## Example (D = 6 and $a_e = 1$ )

$$F\begin{pmatrix}v_{1} \\ v_{2} \\ v_{2} \\ v_{3} \\ v_{4} \\ v_{2} \\ v_{2} \\ v_{2} \\ v_{3} \\ v_{3} \\ v_{3} \\ v_{3} \\ v_{4} \\ v_{2} \\ v_{3} \\ v_{4} \\ v_{3} \\ v_{4} \\ v_{4} \\ v_{4} \\ v_{5} \\ v_$$

$$= \frac{z_{12} - z_{14}}{Q^2} \left[ \ln \frac{Q}{z_3 z_4} \ln \frac{(z_{14} + z_3)(z_{14} + z_4)}{z_{14}(z_{14} + z_3 + z_4)} - \operatorname{Li}_2 \left( \frac{z_3 z_4(z_{14} - z_{12})}{z_{14}Q} \right) \right] \\ + \frac{z_{12} - z_{14}}{Q^2} \operatorname{Li}_2 \left( \frac{z_3 z_4}{Q} \right) + \frac{z_{12}}{Q^2} \ln \frac{z_{14} z_3 z_4}{z_{12}(z_{14} + z_3)(z_{14} + z_4)} - \frac{\ln(z_3 z_4/Q)}{Q(z_{14} + z_3 + z_4)} \right]$$

20 / 24

## Compatibility graph of box-ladders



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- Problem: Find linearly reducible integral representations
  - Schwinger parameters are not always optimal (e.g. for  $K_4$ )

# Summary

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## Thank you!

## Massless $\phi^4$ theory: primitive sixth roots of unity



is not linearly reducible!

Tenth denominator:

$$d_{10} = \alpha_2 \alpha_4^2 \alpha_1 + \alpha_2 \alpha_4^2 \alpha_3 - \alpha_1 \alpha_2 \alpha_3 \alpha_4 + \alpha_2^2 \alpha_4 \alpha_1 + \alpha_2^2 \alpha_4 \alpha_3 - 2\alpha_2 \alpha_3^2 \alpha_4 - \alpha_2^2 \alpha_3^2 - 2\alpha_2^2 \alpha_3 \alpha_1 - 2\alpha_2 \alpha_3^2 \alpha_1 - \alpha_3^2 \alpha_4^2 - 2\alpha_3^2 \alpha_4 \alpha_1 - \alpha_2^2 \alpha_1^2 - 2\alpha_2 \alpha_3 \alpha_1^2 - \alpha_3^2 \alpha_1^2.$$

Change variables:  $\alpha_3 = \frac{\alpha'_3 \alpha_1}{\alpha_1 + \alpha_2 + \alpha_4}$ ,  $\alpha_4 = \alpha'_4 (\alpha_2 + \alpha'_3)$  and  $\alpha_1 = \alpha'_1 \alpha'_4$ ,

$$d_{10}' = (\alpha_2 + \alpha_3')(\alpha_2 + \alpha_2\alpha_4' - \alpha_1')(\alpha_1'\alpha_4' + \alpha_2 + \alpha_2\alpha_4' + \alpha_3'\alpha_4')$$

Final result: not a multiple zeta value, instead MPL at  $e^{i\pi/3}$ 

 $\sqrt{3}\mathcal{P}(P_{7,11})$ 

$$\begin{split} &= \mathsf{Im} \left( \frac{19\,285}{6} \zeta_9 \,\mathsf{Li}_2 - \frac{1029}{2} \zeta_7 \,\mathsf{Li}_4 + 240 \zeta_3^2 \big(9 \,\mathsf{Li}_{2,3} - 7\zeta_3 \,\mathsf{Li}_2\big) \right) - \frac{93\,824}{9675} \pi^3 \zeta_{3,5} \\ &+ \frac{2592}{215} \,\mathsf{Im} \left( 36 \,\mathsf{Li}_{2,2,2,5} + 27 \,\mathsf{Li}_{2,2,3,4} + 9 \,\mathsf{Li}_{2,2,4,3} + 9 \,\mathsf{Li}_{2,3,2,4} + 3 \,\mathsf{Li}_{2,3,3,3} \right. \\ &- 43 \zeta_3 \big(\mathsf{Li}_{2,3,3} + 3 \,\mathsf{Li}_{2,2,4}\big) \Big) - \frac{96\,393\,596\,519\,864\,341\,538\,701\,979}{790\,371\,465\,315\,684\,594\,157\,620\,000} \pi^{11} \\ &+ \frac{216}{14\,755\,731\,798\,995} \,\mathsf{Im} \left( 2\,539\,186\,130\,125\,890\,\mathsf{Li}_8\,\zeta_3 - 1\,269\,593\,065\,062\,945\,\mathsf{Li}_{2,9} \right. \\ &- 413\,965\,317\,054\,502\,\mathsf{Li}_6\,\zeta_5 - \,996\,412\,983\,391\,539\,\mathsf{Li}_{3,8} \\ &- 546\,306\,741\,059\,841\,\mathsf{Li}_{4,7} \, - \,156\,228\,639\,992\,955\,\mathsf{Li}_{5,6} \, \Big) \\ &+ \frac{2592}{10\,945\,435} \pi^2 \,\mathsf{Im} \left( 287\,205\,\mathsf{Li}_{2,7} - 574\,410\,\mathsf{Li}_6\,\zeta_3 + 55\,687\,\mathsf{Li}_{4,5} + 168\,941\,\mathsf{Li}_{3,6} \right. \\ &+ \pi \left( \frac{11\,613\,751}{9030}\,\zeta_5^2 + \frac{267\,067}{602}\,\zeta_{3,7} - \frac{31\,104}{215}\,\mathsf{Re}(3\,\mathsf{Li}_{4,6}\,+10\,\mathsf{Li}_{3,7}) \right) \end{split}$$

Abbreviation:  $\text{Li}_{n_1,...,n_r} := \text{Li}_{n_1,...,n_r}(e^{i\pi/3})$