

HIGGS PRODUCTION: CALCULATING AT N3LO

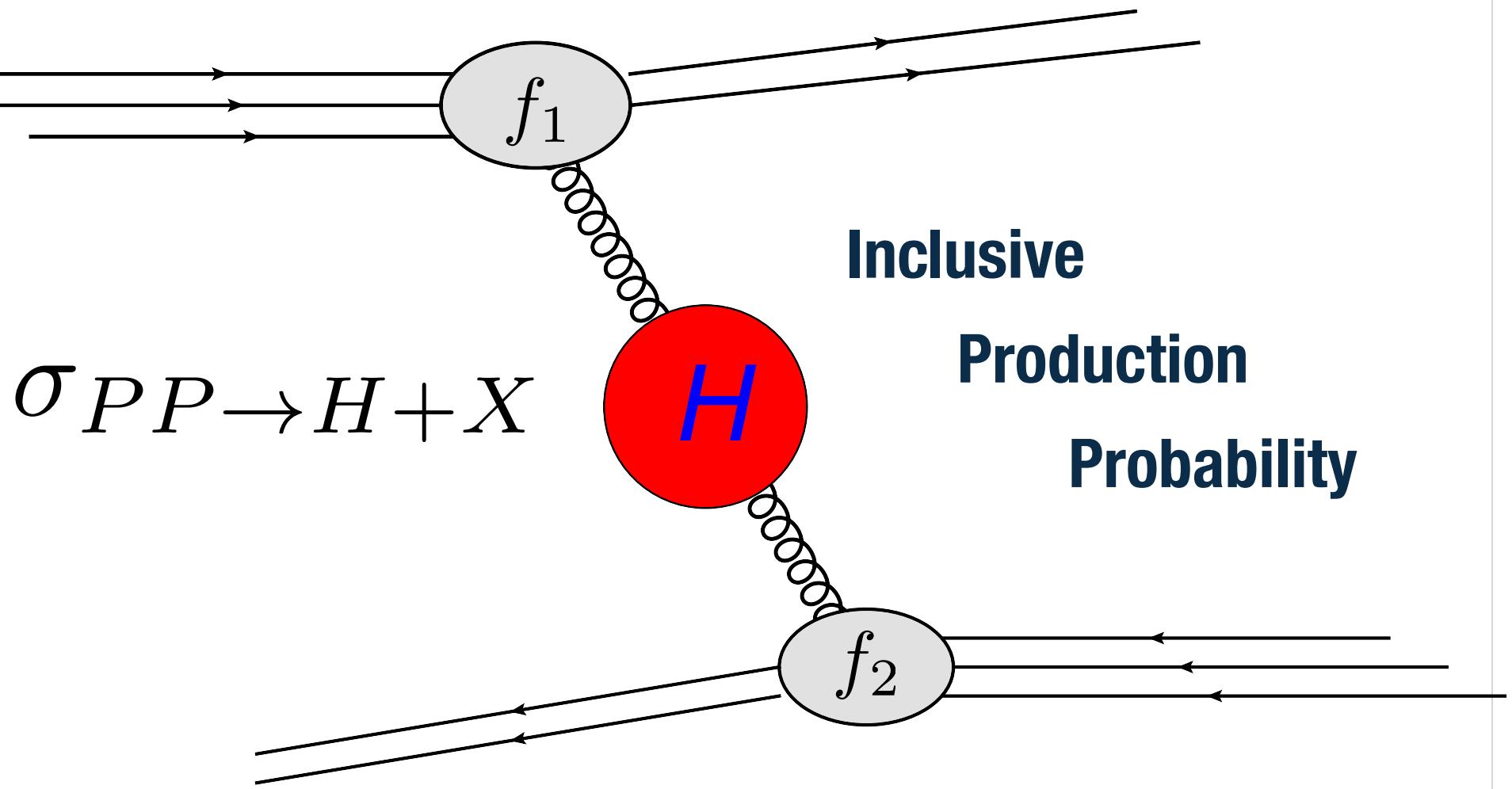
Bernhard Mistlberger

ETH Zurich

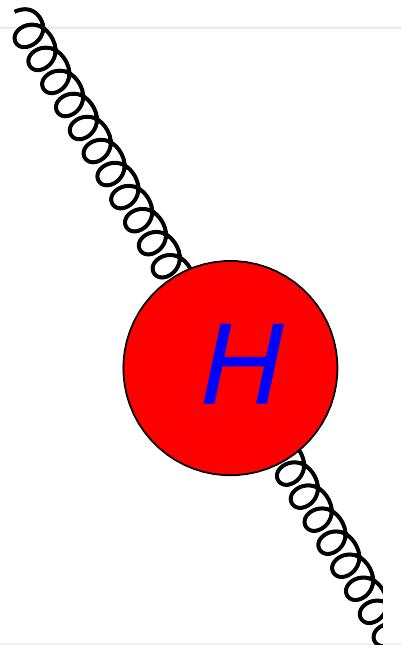
The N3LO Team:

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Elisabetta Furlan, Franz Herzog, Thomas Gehrmann,
Achilleas Lazopoulos, BM

HIGGS PRODUCTION



HIGGS PRODUCTION

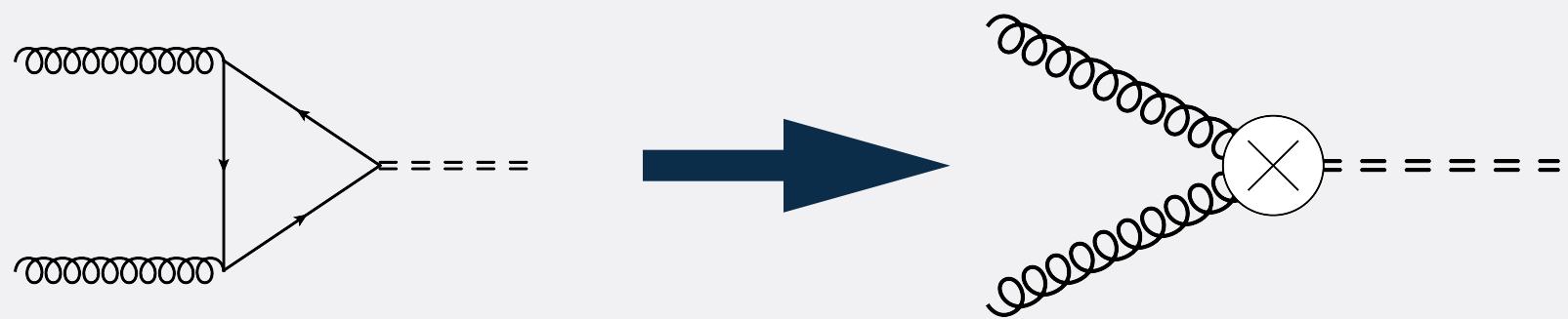


Calculate from perturbative QFT+PDF

$$\sigma_{PP \rightarrow H + X} = \int dx_1 dx_2 f_1(x_1) f_2(x_2) \hat{\sigma}(x_1 x_2)$$

QCD dominates at the LHC

Gluon-Fusion

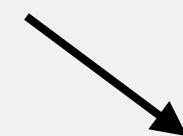


Heavy Top Approximation

GLUON FUSION

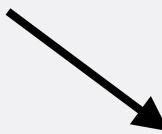
1978 **L0**

$$\hat{\sigma}^{8TeV} = 9.6pb$$
$$\pm 25\%$$



1991 **NLO**

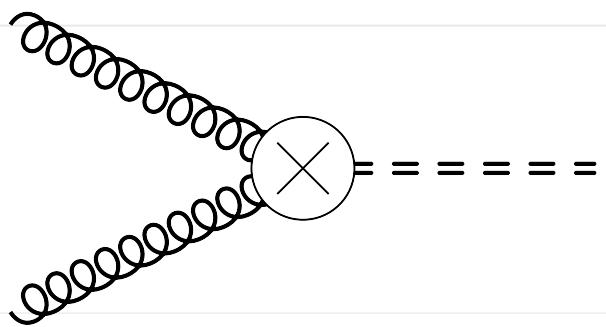
$$\hat{\sigma}^{8TeV} = 16.7pb$$
$$\pm 20\%$$



2002

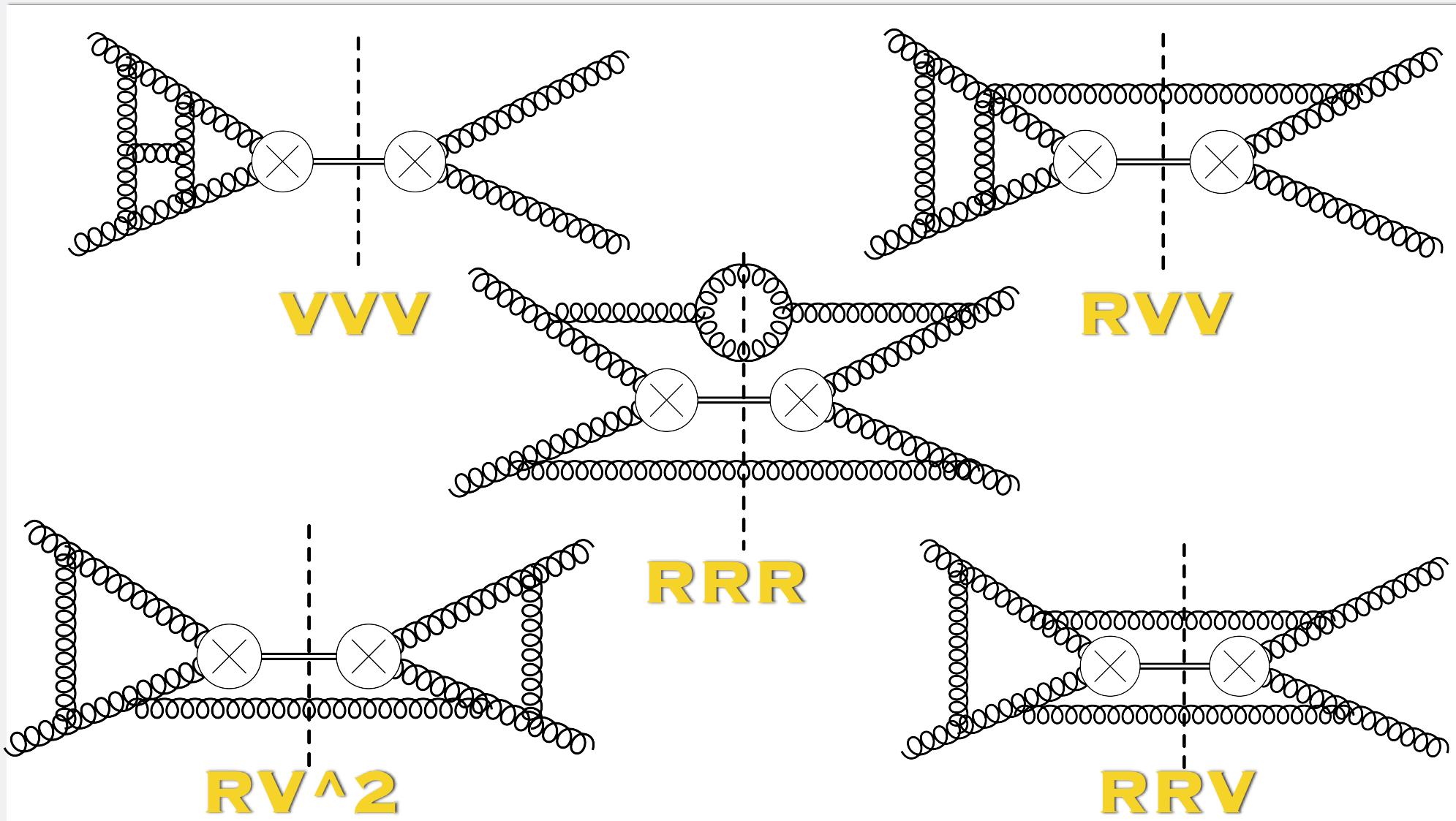
NNLO

$$\hat{\sigma}^{8TeV} = 19.6pb$$
$$\pm 9\%$$

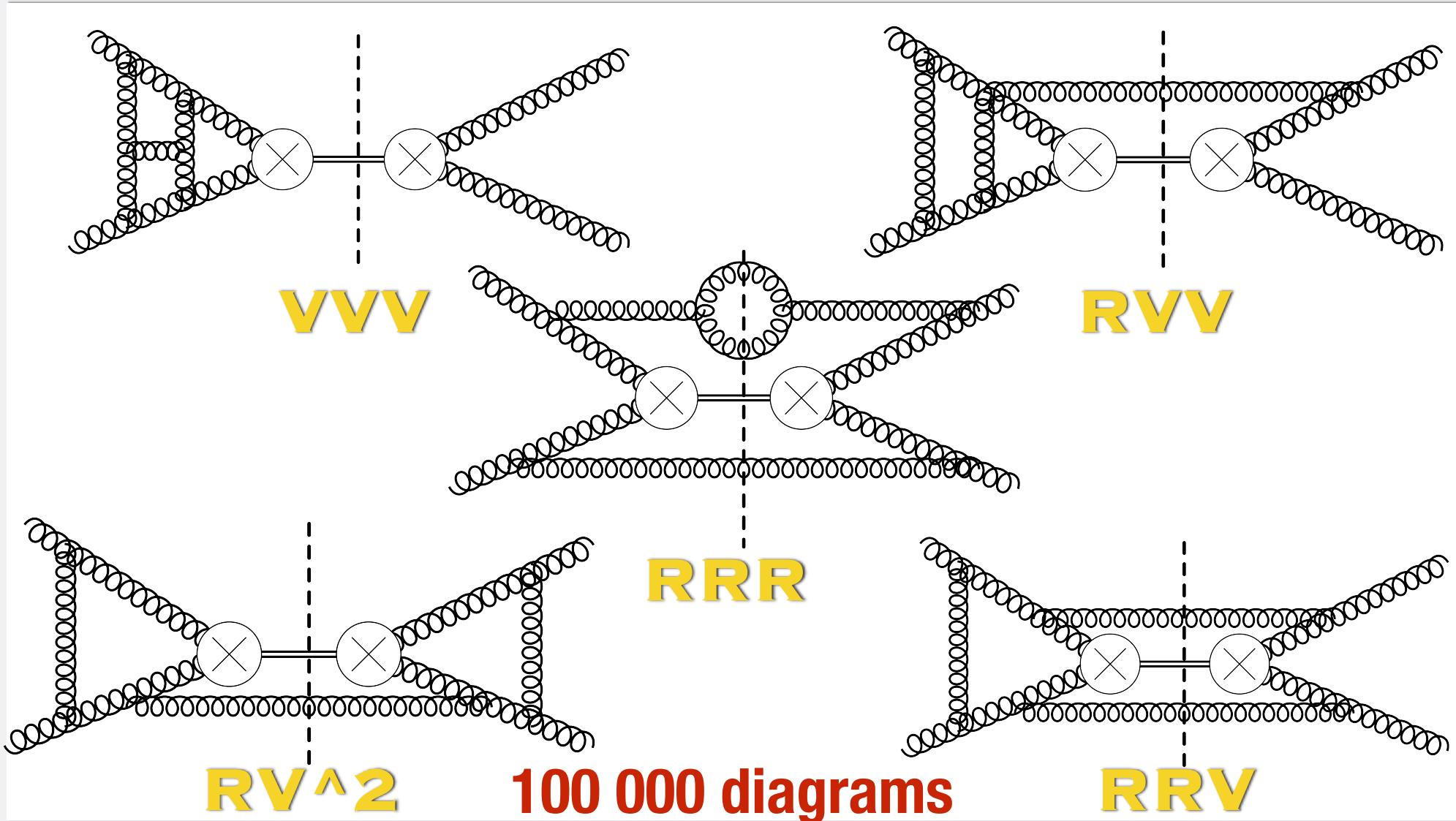


2015 **N3LO**

FEYNMAN DIAGRAMS

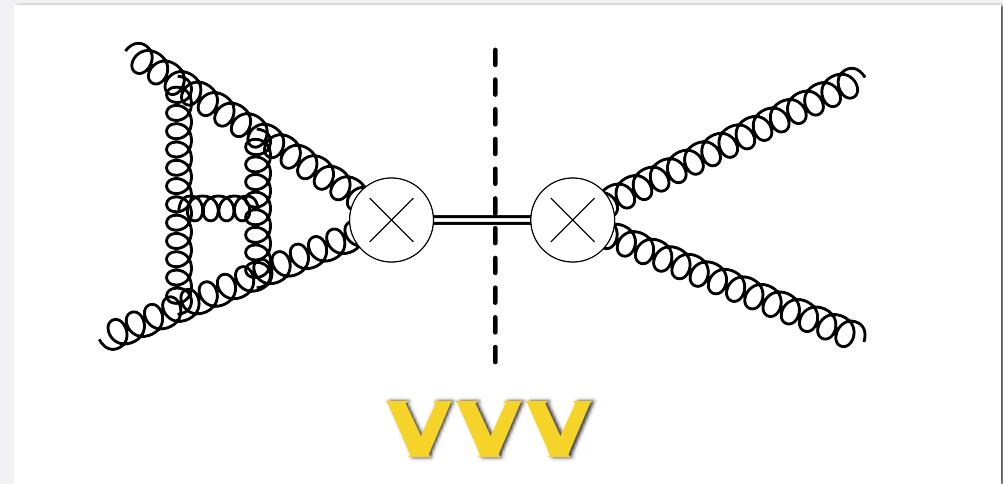


FEYNMAN DIAGRAMS



THE INTEGRAND

Loop Integrals



$$I = \int \frac{d^d p_h}{(2\pi)^d} \frac{d^d p_4}{(2\pi)^d} \frac{d^d p_5}{(2\pi)^d} \frac{d^d p_6}{(2\pi)^d} \delta_+(p_h^2 - m_h^2) \delta^d(p_1 + p_2 - p_h)$$

$$\times |\mathcal{M}(p_i)|^2$$



VVV Integrand

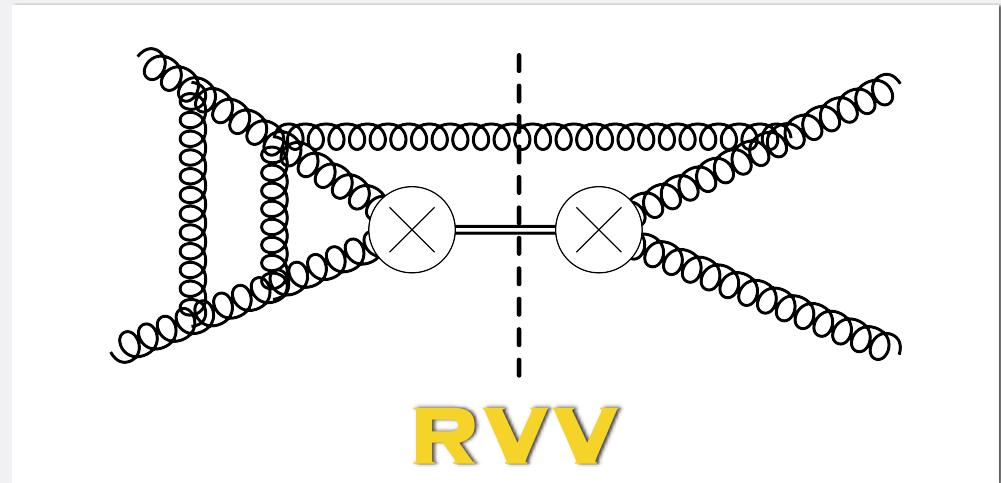
THE INTEGRAND

One parton in
the final state

Loop Integrals

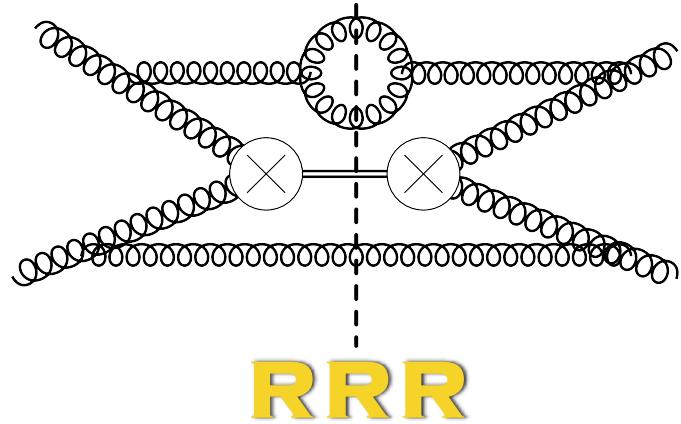
$$I = \int \frac{d^d p_h}{(2\pi)^d} \frac{d^d p_4}{(2\pi)^d} \frac{d^d p_5}{(2\pi)^d} \frac{d^d p_6}{(2\pi)^d} \delta_+(p_h^2 - m_h^2) \delta^d(p_1 + p_2 - p_h - p_4) \\ \times |\mathcal{M}(p_i)|^2 \times \delta_+(p_4^2)$$

RVV Integrand



THE INTEGRAND

No loop left!



$$I = \int \frac{d^d p_h}{(2\pi)^d} \frac{d^d p_4}{(2\pi)^d} \frac{d^d p_5}{(2\pi)^d} \frac{d^d p_6}{(2\pi)^d} \delta_+(p_h^2 - m_h^2) \delta^d(p_1 + p_2 - p_h - p_4 - p_5 - p_6)$$

$$\times |\mathcal{M}(p_i)|^2 \times \delta_+(p_4^2) \times \delta_+(p_5^2) \times \delta_+(p_6^2)$$



RRR Integrand

THE INTEGRAND

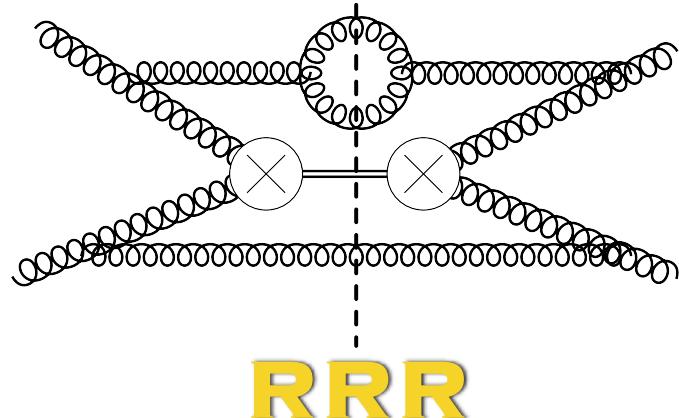
No loop left!

Integrals depend only on

s

and

$$z = \frac{m_h^2}{s}$$



$$I = \int \frac{d^d p_h}{(2\pi)^d} \frac{d^d p_4}{(2\pi)^d} \frac{d^d p_5}{(2\pi)^d} \frac{d^d p_6}{(2\pi)^d} \delta_+(p_h^2 - m_h^2) \delta^d(p_1 + p_2 - p_h - p_4 - p_5 - p_6)$$
$$\times |\mathcal{M}(p_i)|^2 \times \delta_+(p_4^2) \times \delta_+(p_5^2) \times \delta_+(p_6^2)$$



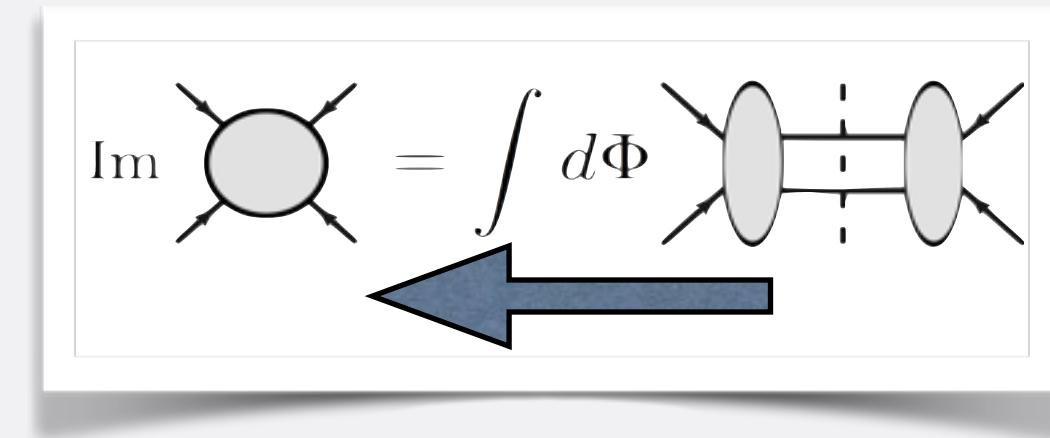
RRR Integrand

Exploit similarity!

REDUCING COMPLEXITY

REVERSE UNITARITY

Optical theorem



Cutkosky's Rule

$$\delta^+(p^2) \rightarrow \left[\frac{1}{p^2} \right]_c \sim \frac{1}{p^2 + i\epsilon} - \frac{1}{p^2 - i\epsilon}$$

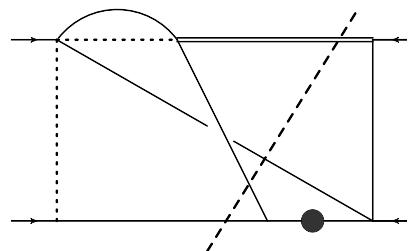
Treat cuts almost as ordinary propagators!

REDUCING COMPLEXITY

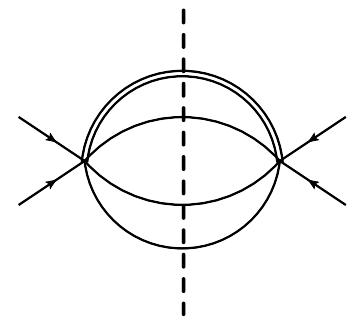
Use “Integration-By-Part” (IBP) Identities

$$\int \frac{d^d p_4}{(2\pi)^d} \frac{\partial}{\partial p_4^\mu} \left(p_4^\mu \left[\frac{1}{p_4^2} \right]_c |\mathcal{M}(p_4, \dots)|^2 \dots \right) = 0$$

→ “Master Integrals”



$$= - \frac{(\epsilon - 1)(2\epsilon - 1)(3\epsilon - 2)(3\epsilon - 1)(6\epsilon - 5)(6\epsilon - 1)}{\epsilon^4(\epsilon + 1)(2\epsilon - 3)}$$



REDUCING COMPLEXITY

Use “Integration-By-Part” Identities



“Master Integrals”

A Feynman diagram on the left shows a loop with a vertical line entering from the left and a horizontal line exiting to the right. A black dot is at the bottom vertex. A dashed line connects the top vertex to the right. A dotted line connects the top vertex to the left. A solid line connects the left vertex to the bottom vertex. An equals sign follows the diagram. To the right is a complex rational function:

$$= -\frac{(\epsilon - 1)(2\epsilon - 1)(3\epsilon - 2)(3\epsilon - 1)(6\epsilon - 5)(6\epsilon - 1)}{\epsilon^4(\epsilon + 1)(2\epsilon - 3)}$$

A much simpler Feynman diagram on the right consists of two nested circles with arrows indicating direction. A vertical dashed line passes through the center of both circles.

Does that simplify life?

before

NNLO

50 000

after

27

REDUCING COMPLEXITY

Use “Integration-By-Part” Identities

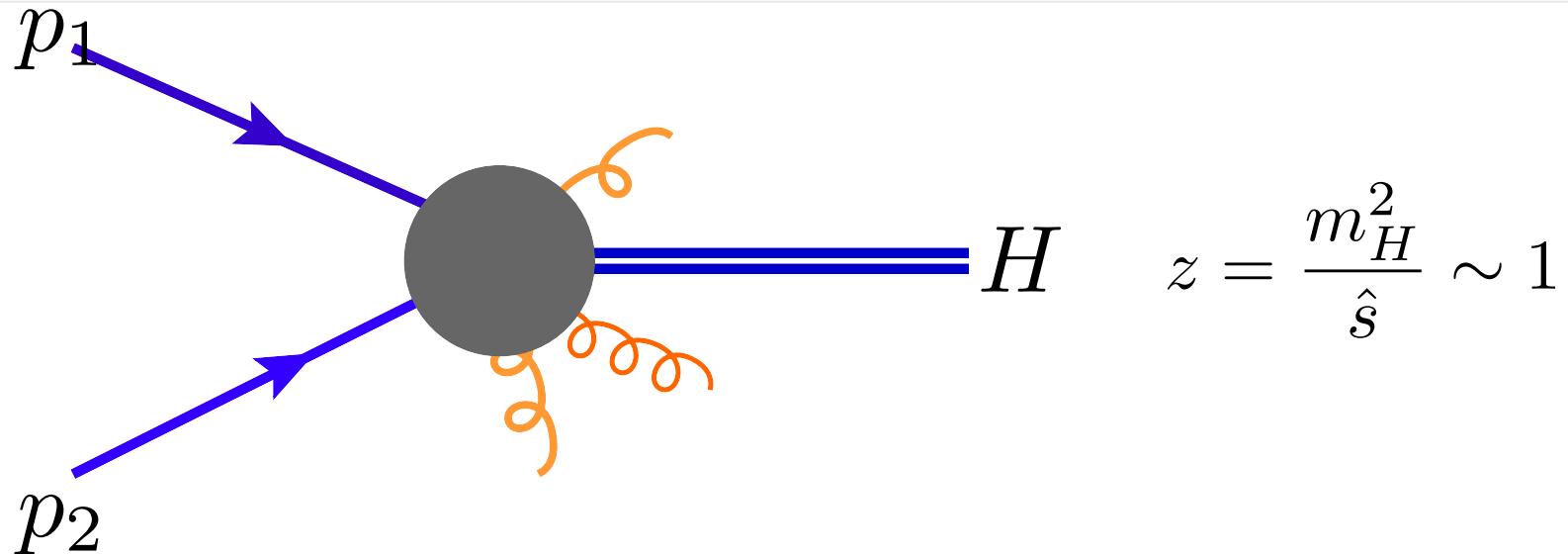
→ “Master Integrals”

A Feynman diagram consisting of a rectangle with a semi-circular arc on top. A dashed line from the top-right corner divides the rectangle into two triangles. A solid black dot is at the bottom-right corner. To the right of the diagram is an equals sign followed by a complex rational function of ϵ :
$$= -\frac{(\epsilon - 1)(2\epsilon - 1)(3\epsilon - 2)(3\epsilon - 1)(6\epsilon - 5)(6\epsilon - 1)}{\epsilon^4(\epsilon + 1)(2\epsilon - 3)}$$

To the right of the equation is a smaller diagram showing a circle with a vertical dashed line through its center, representing a master integral.

Does that simplify life?	before	after
NNLO	50 000	27
N3LO	517 531 178	1028

THRESHOLD EXPANSION

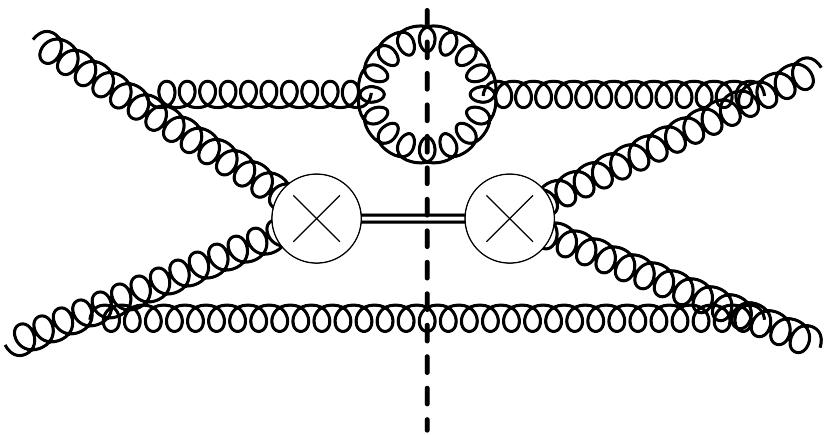


Expand around the production threshold of the Higgs

$$\bar{z} = 1 - z \quad \rightarrow \quad \hat{\sigma}(\bar{z}) = \sigma^{SV} + \sigma^{(0)} + \bar{z}\sigma^{(1)} + \dots$$

MASTER INTEGRALS

Method #1: **Momentum Space Expansion**



Rescale final state momenta

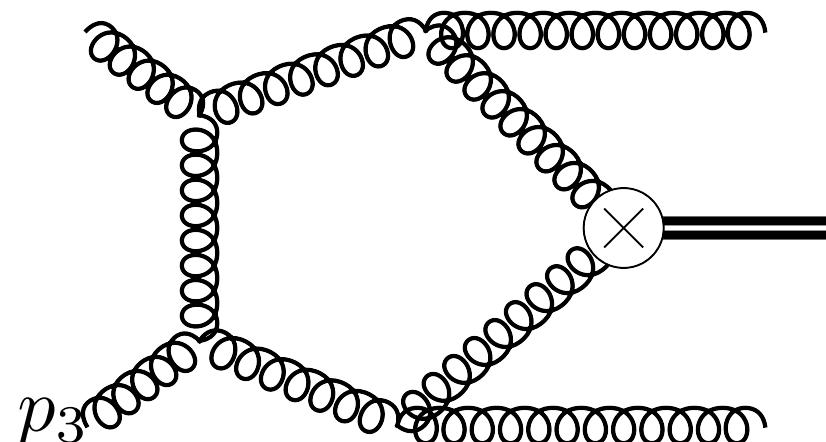
$$p_f \rightarrow (1 - z)p_f$$

Expand integrand and measure

Ready to compute RRR!

MASTER INTEGRALS

Method #1: Momentum Space Expansion



Not so easy with loops: RRV

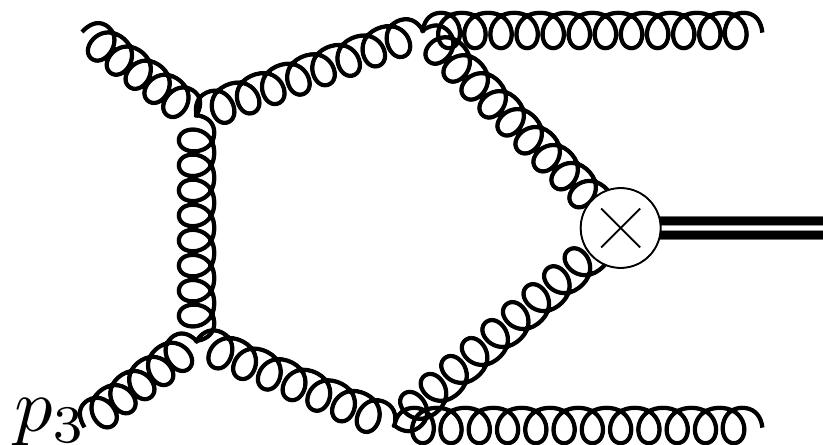
**Loop momentum
not constrained**

Split in regions

Soft	Coll 1	Coll 2	Hard
$k \rightarrow \bar{z}k$	$k \rightarrow k p_1$	$k \rightarrow k p_2$	k

MASTER INTEGRALS

Method #1: Momentum Space Expansion



parametrize loop momentum

$$p_l \rightarrow p_1\alpha + p_2\beta + p_{l,\perp}$$

Hard

$$\alpha \rightarrow \alpha, \quad \beta \rightarrow \beta, \quad p_{l,\perp}^2 \rightarrow p_{l,\perp}^2$$

Coll 1

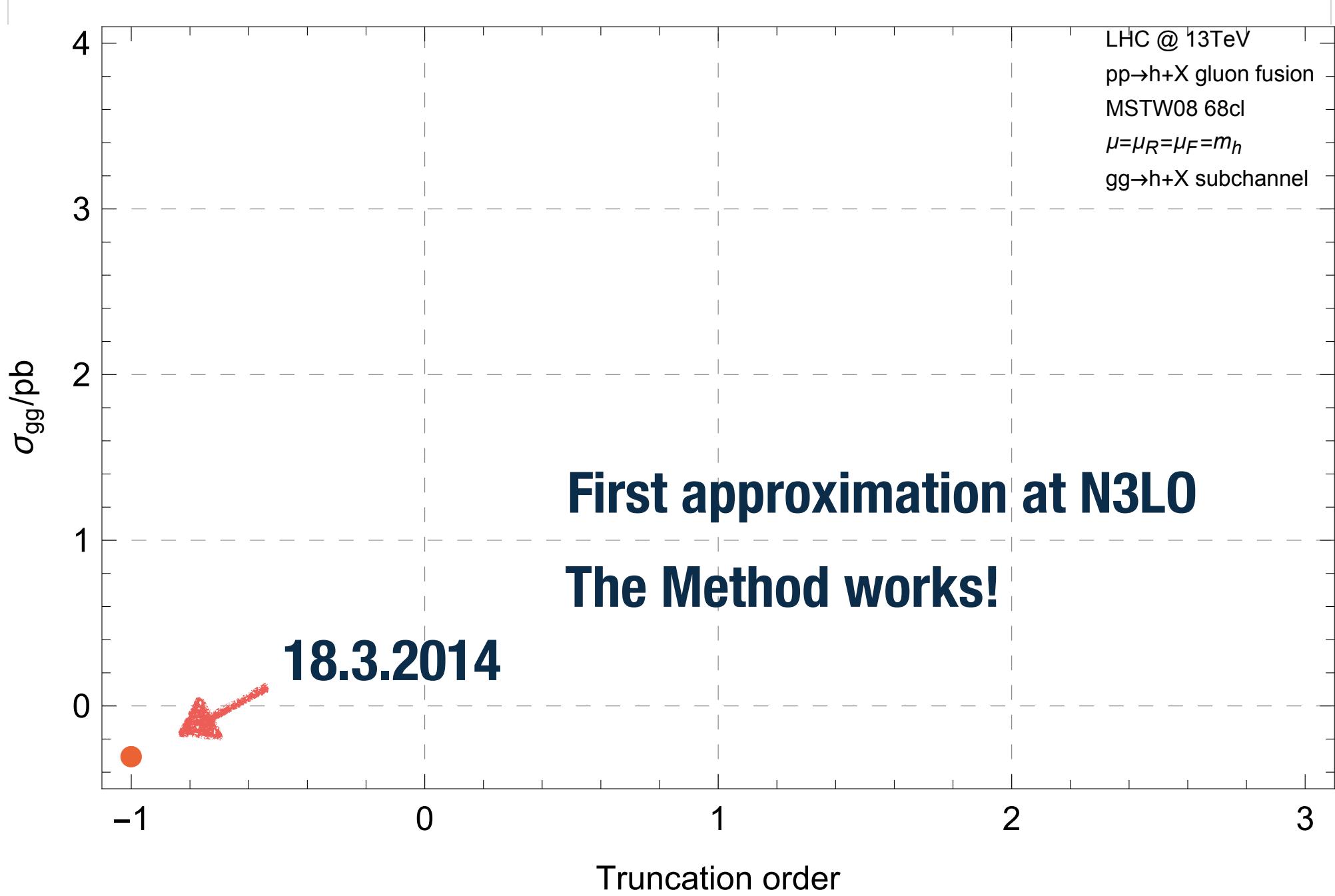
$$\alpha \rightarrow \alpha, \quad \beta \rightarrow (1-z)\beta, \quad p_{l,\perp}^2 \rightarrow (1-z)p_{l,\perp}^2$$

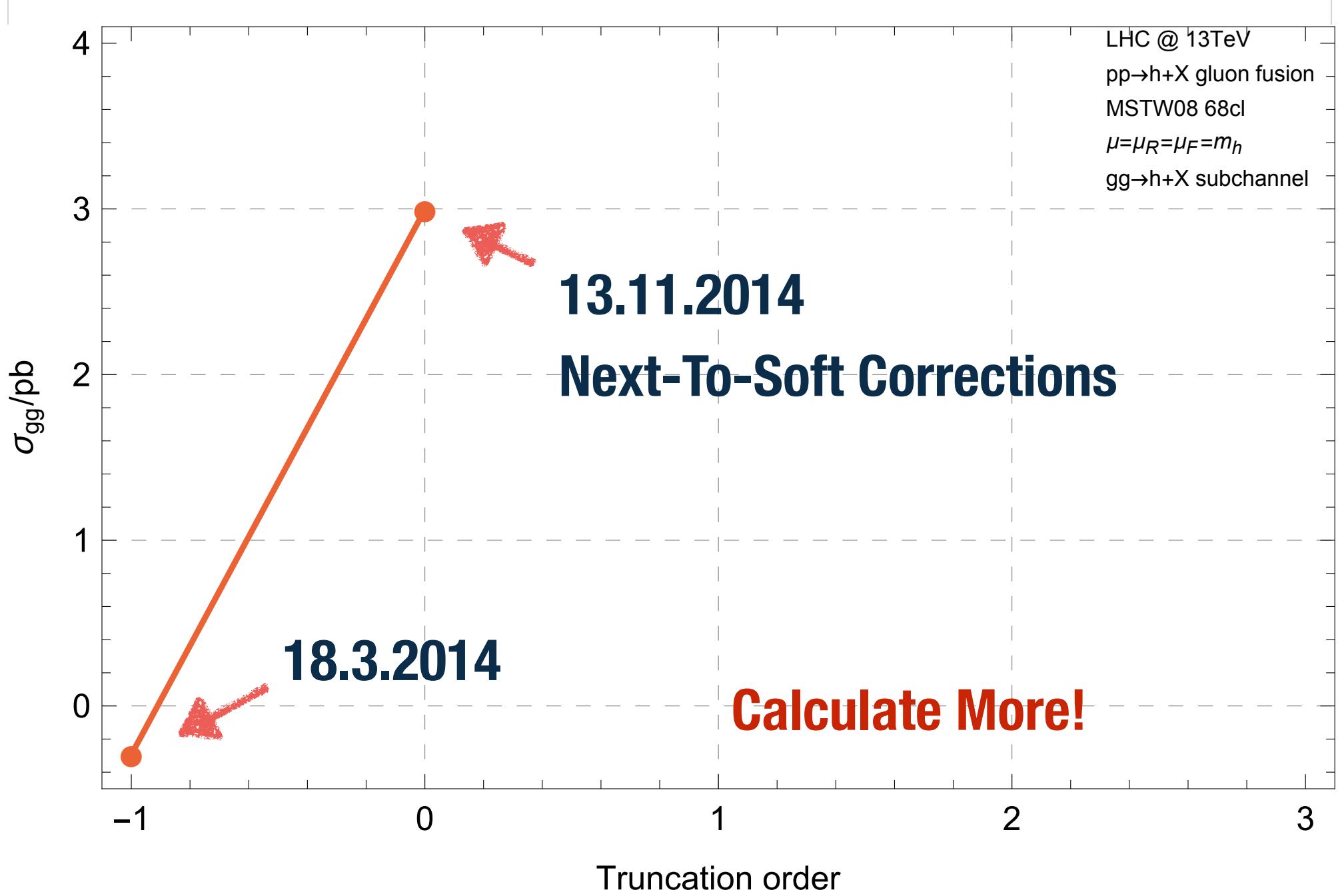
Coll 2

$$\alpha \rightarrow (1-z)\alpha, \quad \beta \rightarrow \beta, \quad p_{l,\perp}^2 \rightarrow (1-z)p_{l,\perp}^2$$

Soft

$$\alpha \rightarrow (1-z)\alpha, \quad \beta \rightarrow (1-z)\beta, \quad p_{l,\perp}^2 \rightarrow (1-z)p_{l,\perp}^2$$





MASTER INTEGRALS

Method #2: **Differential Equations**

$$z = \frac{m_h^2}{s}$$

Only two variables: s and z

Dependence on s trivial: „Energy - dimension“

Integrand depends on z
only via Higgs on-shell constraint

$$\frac{\partial}{\partial z} \delta(p_h^2 - sz) \rightarrow s \left(\frac{1}{p_h^2 - sz} \right)_c^2$$

Relate the differential with IBP identities to Master-Integrals

$$\frac{\partial}{\partial z} \vec{M} = A(z) \vec{M}$$

MASTER INTEGRALS

Method #2: **Differential Equations**

$$\frac{\partial}{\partial z} \vec{M} = A(\epsilon, z) \vec{M}$$

- **System of coupled differential equations**
- **Coefficient Matrix A is not constant**

One solution:

Find a way to decouple order by order in ϵ

$$\frac{\partial}{\partial z} \vec{M} = \epsilon A(\epsilon, z) \vec{M}$$

MASTER INTEGRALS

Solve

$$\frac{\partial}{\partial z} \vec{M} = A(\epsilon, z) \vec{M}$$

with a generalised power series Ansatz

$$M_i = \sum_j \sum_{k=2}^6 c_{ijk} (1 - z)^{(j-k\epsilon)}$$

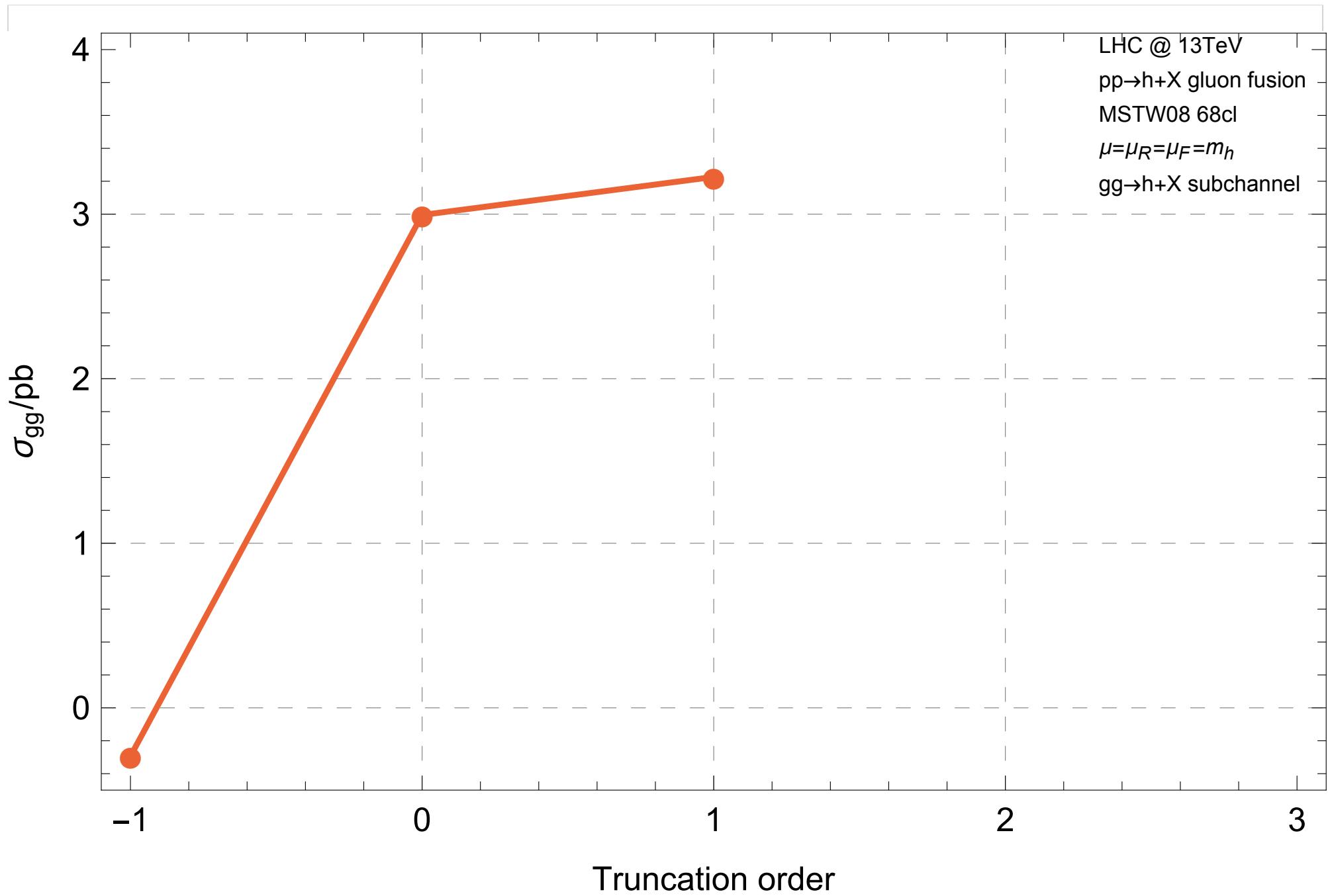
1028 Differential Equations
need 1028 Boundary Conditions!

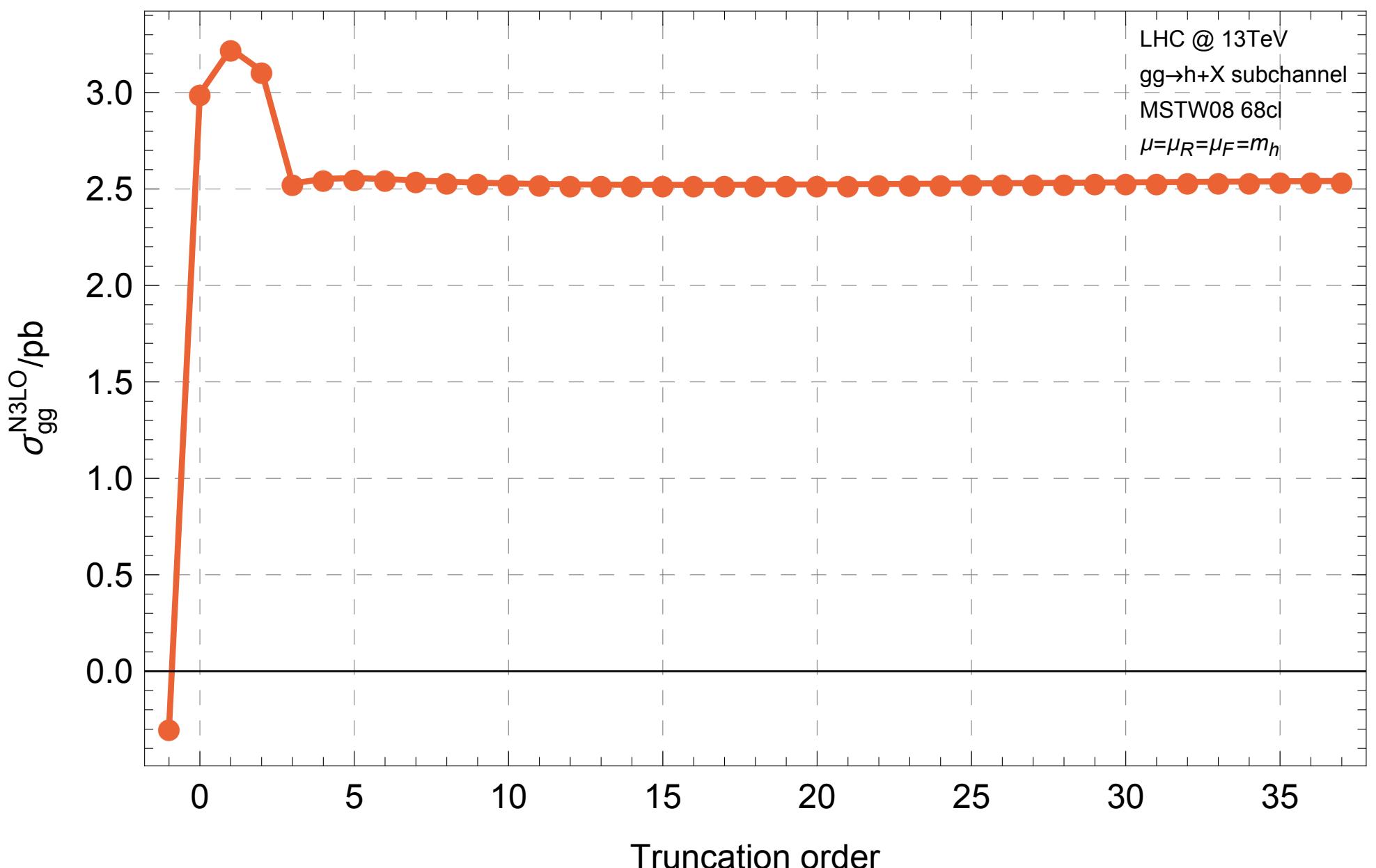
Combine with Expansion-by-Regions

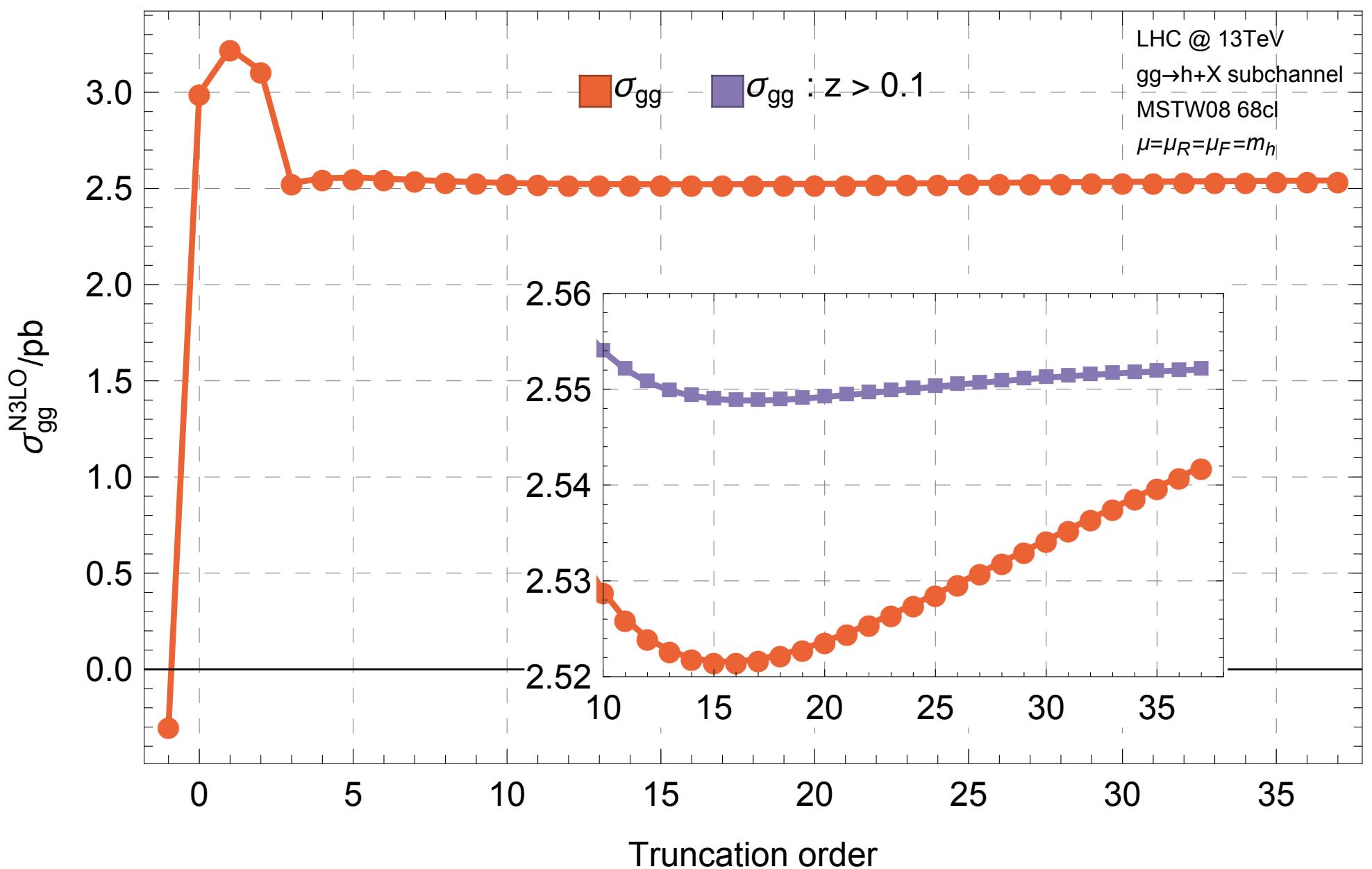
1028



78







THE METHODS

- Technology to handle hundreds of thousands of Feynman diagrams
- Improved algorithms to reduce the complexity of integrands to a manageable size
- Techniques to perform integrand and integral expansions
- Novel methods of computing Feynman Integrals

Ready to predict at N3LO!

CONCLUSIONS

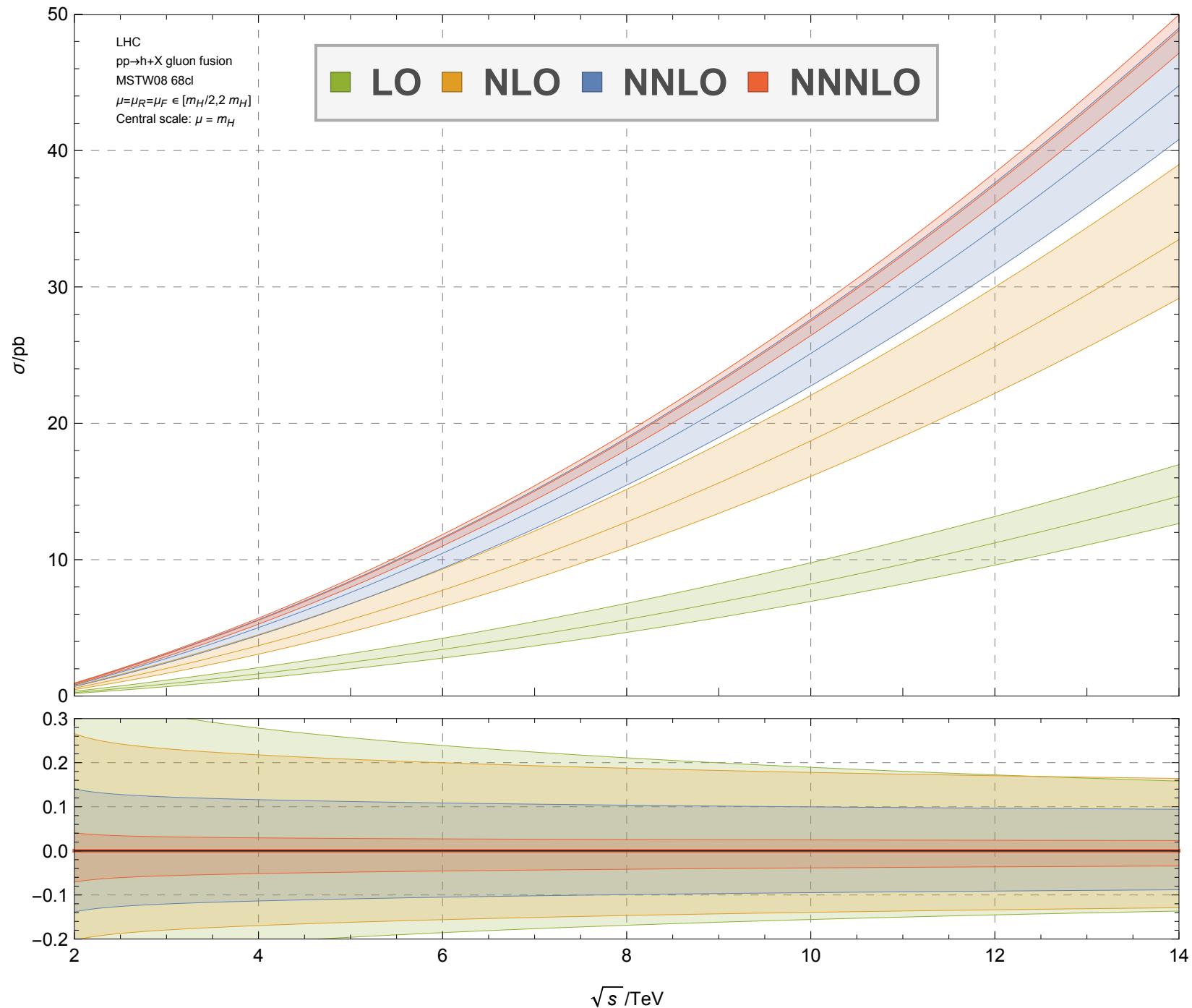
First complete calculation of an observable at N3LO

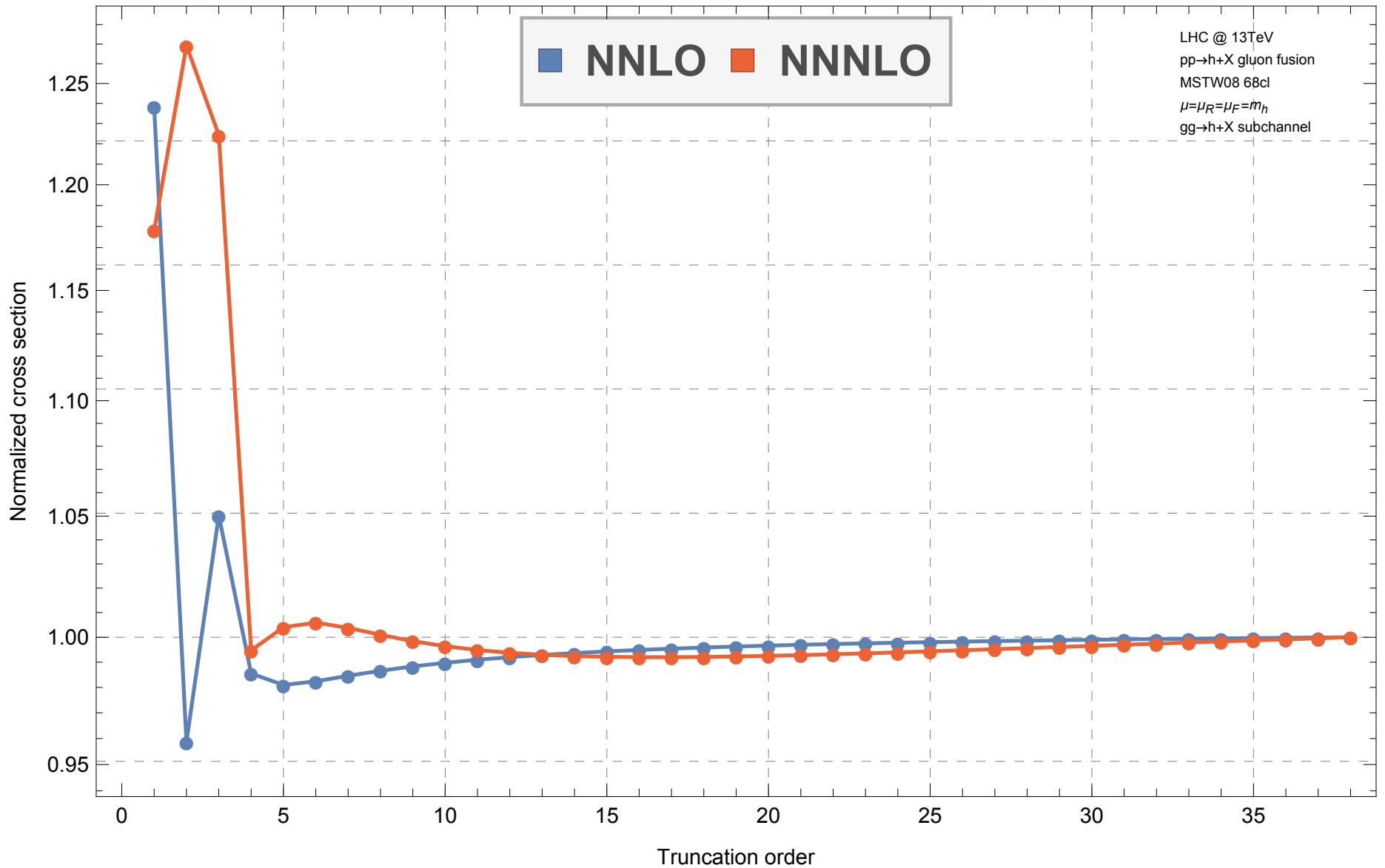
30 orders in the threshold expansion for Gluon Fusion

Drastic reduction of scale dependence

Direct impact on LHC phenomenology

This is the dawn of the age of N3LO precision





PARTON LUMINOSITY

$$\mathcal{L}_{12}(z) = \int_{\tau/z}^1 dx_1 f_1(x_1) f_2\left(\frac{\tau}{zx_1}\right)$$

