

# NNLO Corrections to Dijet Production

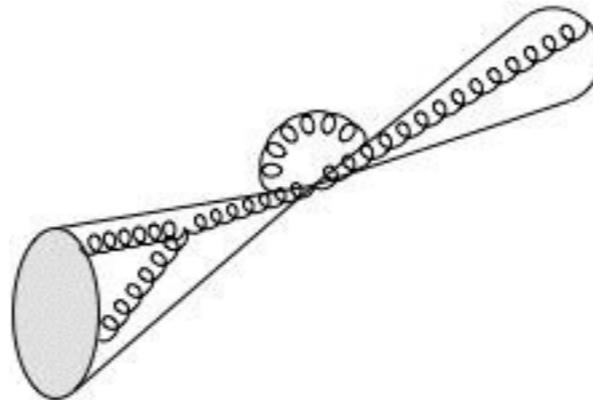
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in direct collaboration with:

Nigel Glover (IPPP), João Pires (Milan) and Steven Wells (IPPP)



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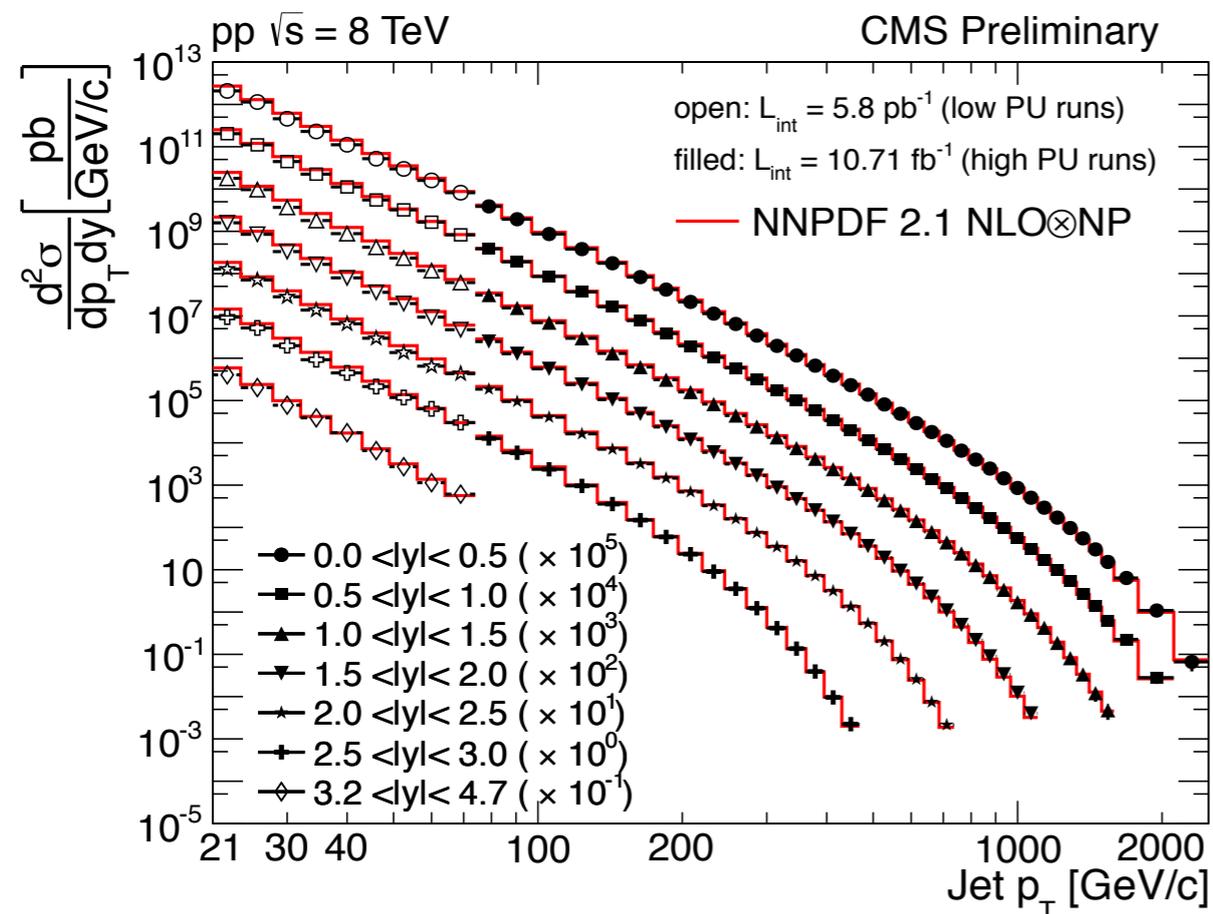
**MC@NNLO**

# Jets at the LHC

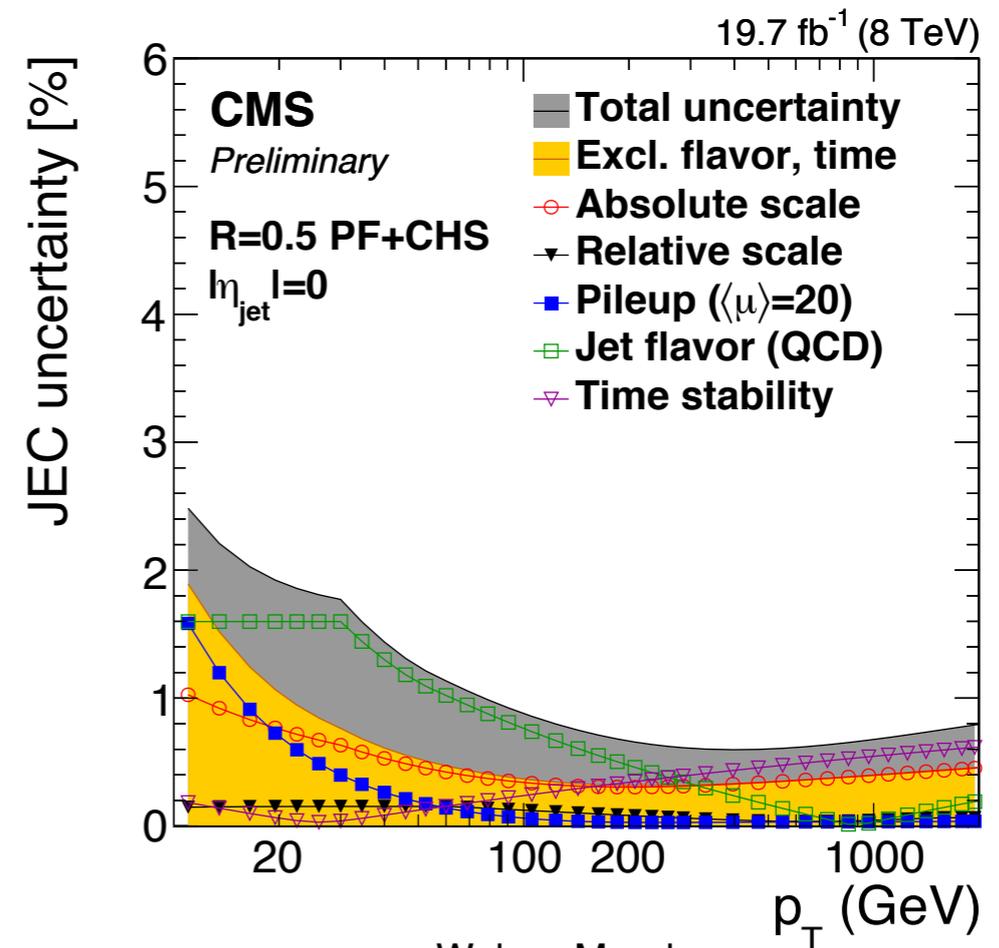
Ubiquitous and accurately measured at the LHC

- ~1% JES corresponds to <10% uncertainty on single inclusive x-sec

Provides a rigorous test of QCD across a huge range of kinematic variables



CMS-PAS-FSQ-12-031

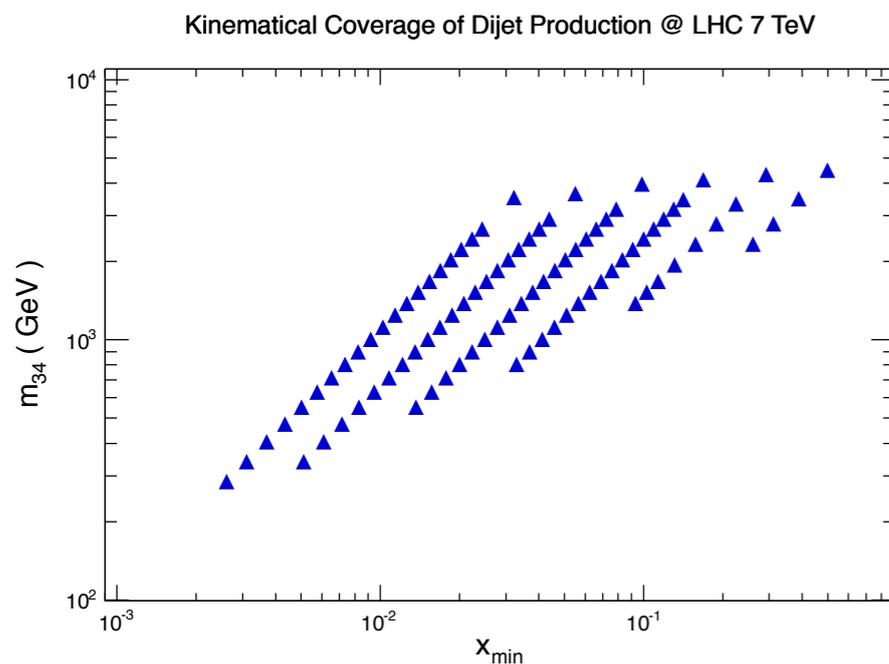


Weber, Monday

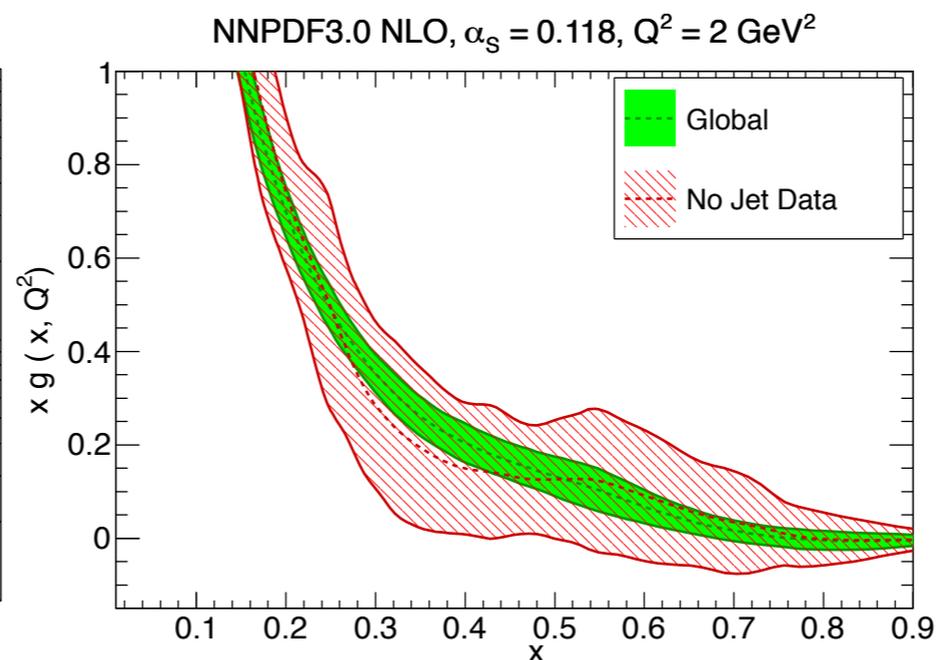
# Jets and PDFs

LHC is mainly a gluon collider but gluon PDF is not well known:

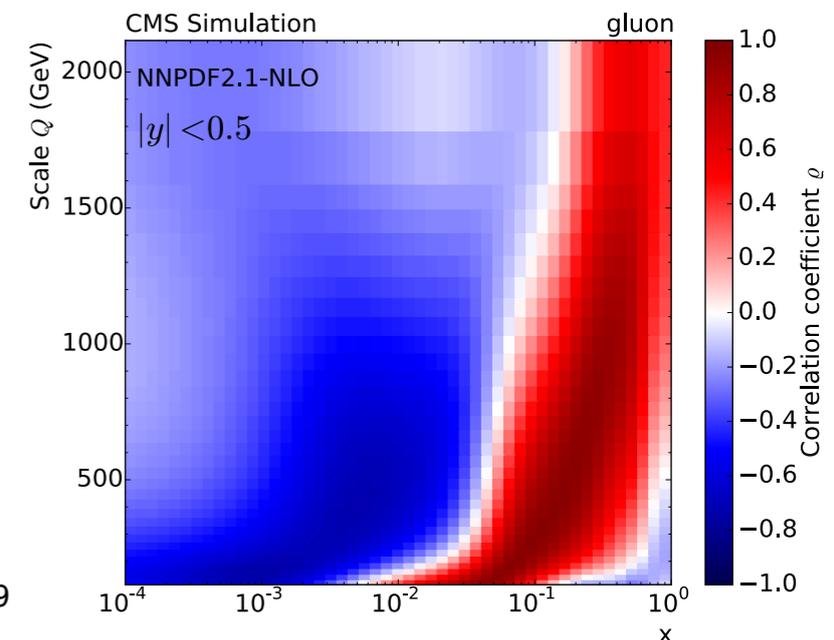
- LHC jets probe a wide range of  $x$
- gluon PDF directly sensitive to jet data, especially at large  $x$
- would like to consistently include NNLO jet data in NNLO PDF fits without using kinematically limited approximations



Rojo hep-ph [1410.7728]



NNPDF collaboration hep-ph [1410.8849]

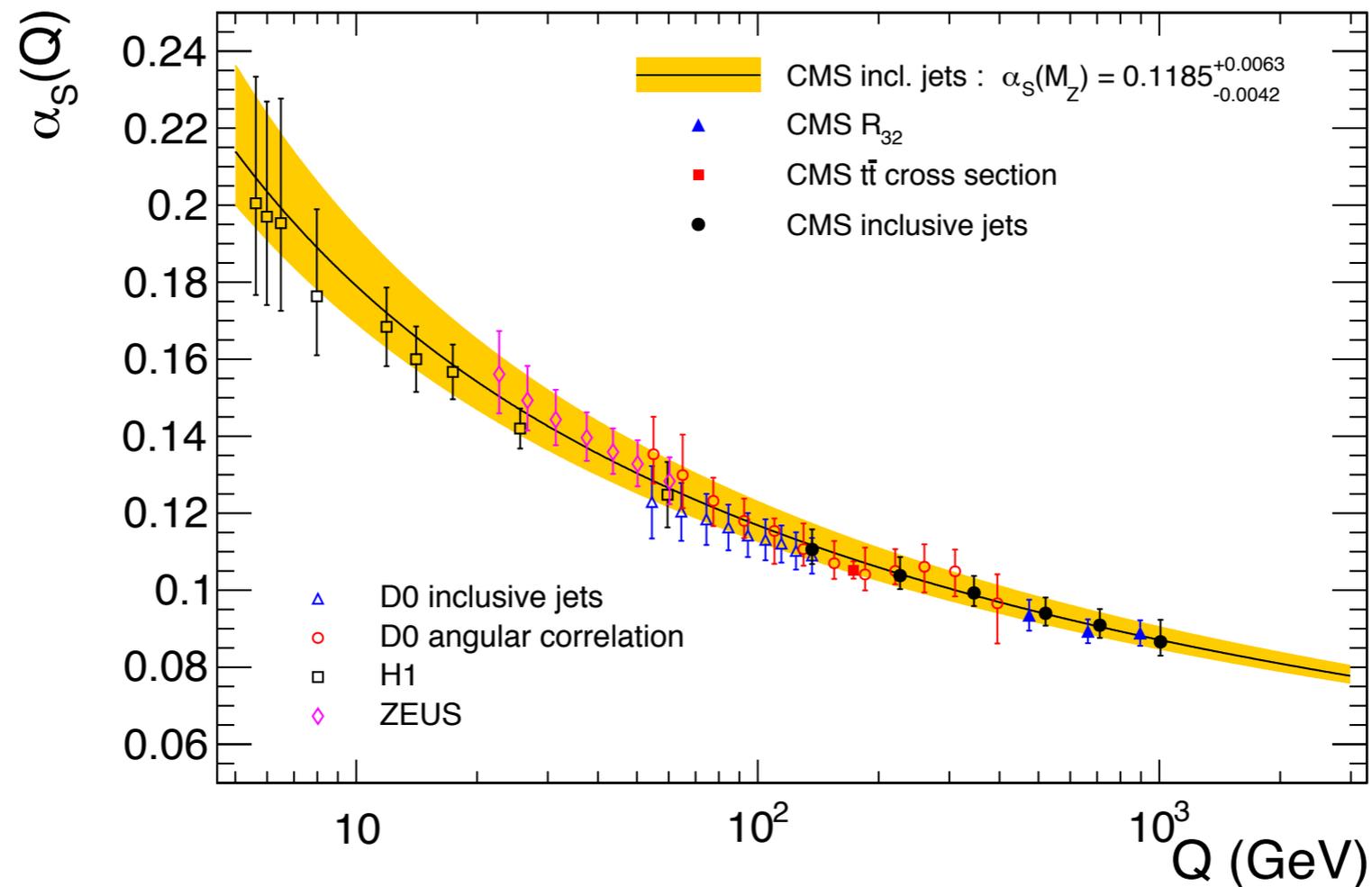


Rojo hep-ph [1410.7728]

# Jets and $\alpha_s$

Can use the single inclusive jet cross section to determine [CMS-PAS-SMP-12-028]:

- $\alpha_s(M_Z)$  and running coupling from single experiment

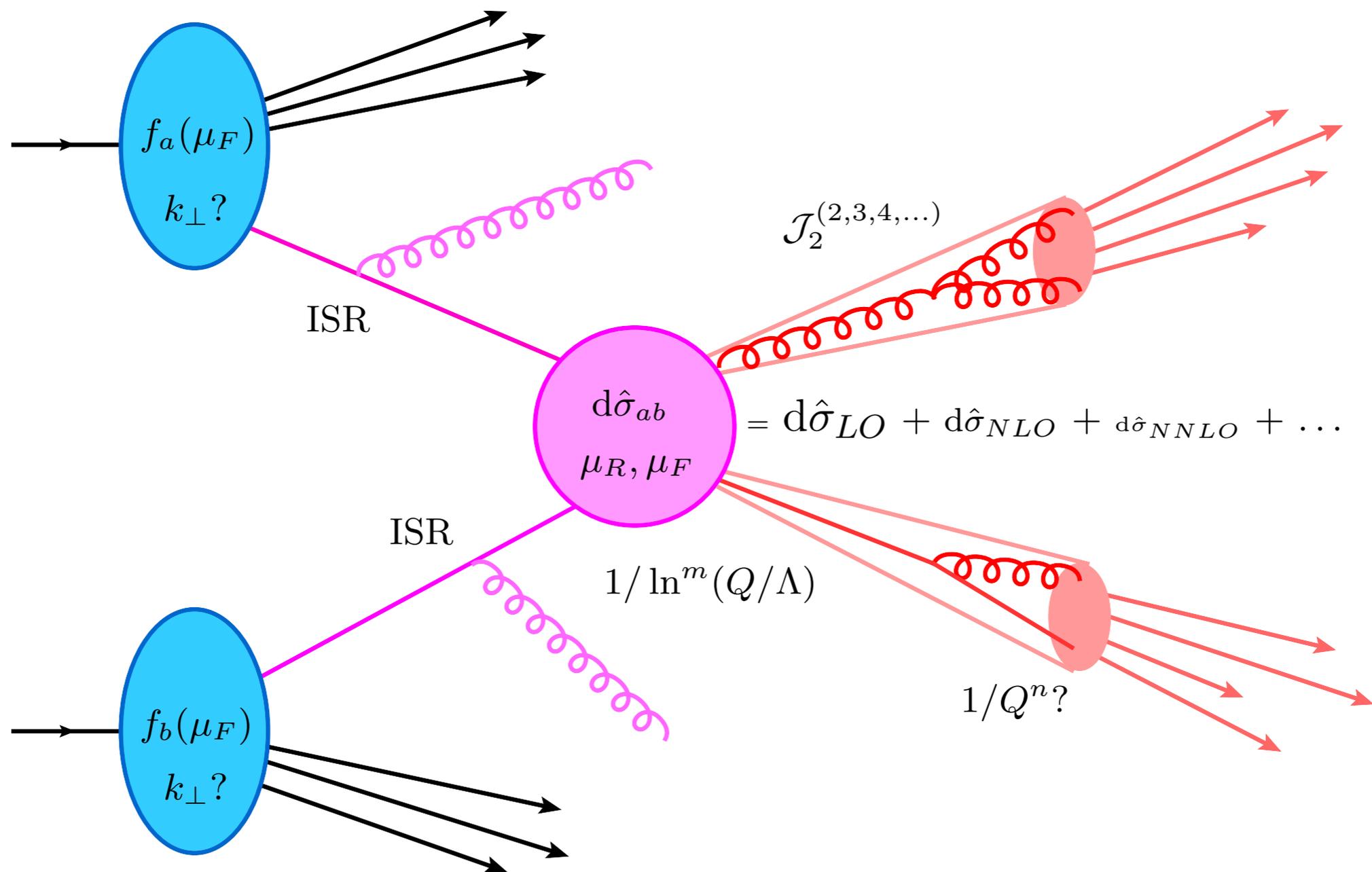


- very satisfying test of QCD and the LHC

CMS hep-ex [1410.6765]

- model independent probe of new physics

# Why NNLO?

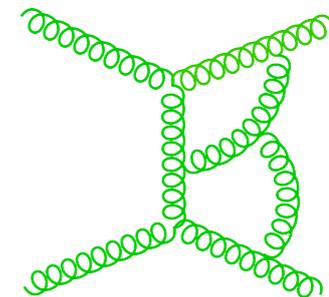
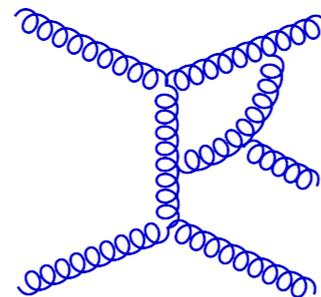
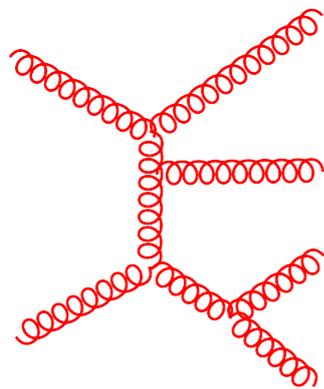


# NNLO Subtraction

Unphysical intermediate quantities are divergent

- need to regulate with RR, RV and VV subtraction terms

$$\begin{aligned}d\sigma_{ab,NNLO} &= \int_{\Phi_{m+2}} d\sigma_{ab,NNLO}^{RR} \\ &+ \int_{\Phi_{m+1}} d\sigma_{ab,NNLO}^{RV} + d\sigma_{ab,NNLO}^{MF,1} \\ &+ \int_{\Phi_m} d\sigma_{ab,NNLO}^{VV} + d\sigma_{ab,NNLO}^{MF,2}\end{aligned}$$

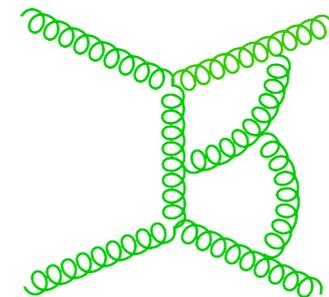
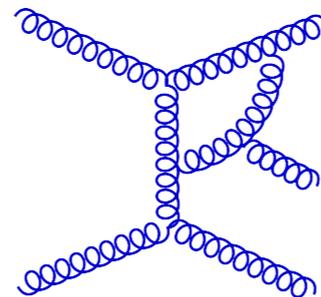
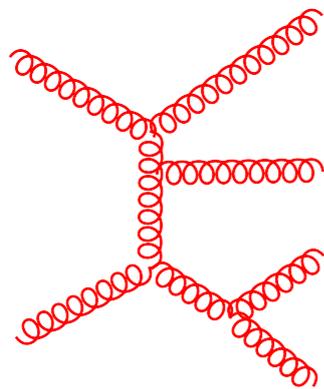


# NNLO Subtraction

Unphysical intermediate quantities are divergent

- need to regulate with RR, RV and VV subtraction terms

$$\begin{aligned} d\sigma_{ab,NNLO} = & \int_{\Phi_{m+2}} \left[ d\sigma_{ab,NNLO}^{RR} - d\sigma_{ab,NNLO}^S \right] \\ & + \int_{\Phi_{m+1}} \left[ d\sigma_{ab,NNLO}^{RV} - d\sigma_{ab,NNLO}^T \right] \\ & + \int_{\Phi_m} \left[ d\sigma_{ab,NNLO}^{VV} - d\sigma_{ab,NNLO}^U \right] \end{aligned}$$



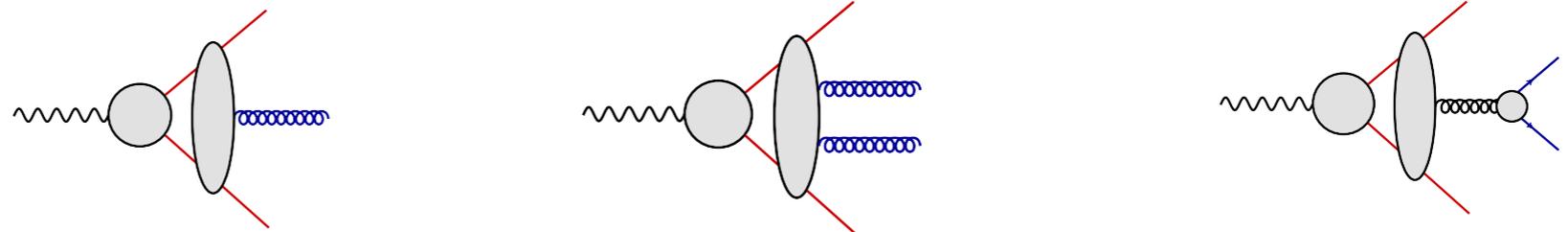
# Antenna Subtraction

Antenna functions built from matrix elements:

$$X_3^0(i, j, k) \sim \frac{|\mathcal{M}_3^0(i, j, k)|^2}{|\mathcal{M}_2^0(I, K)|^2}, \quad X_4^0(i, j, k, l) \sim \frac{|\mathcal{M}_4^0(i, j, k, l)|^2}{|\mathcal{M}_2^0(I, L)|^2}$$

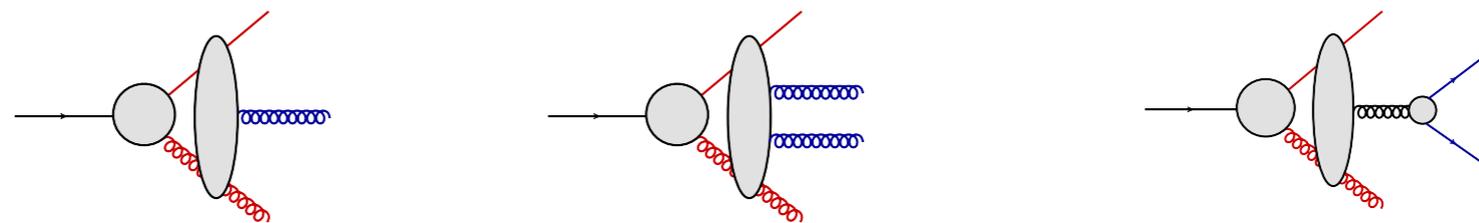
Quark-antiquark:

$$\gamma^* \rightarrow q\bar{q} + \dots$$



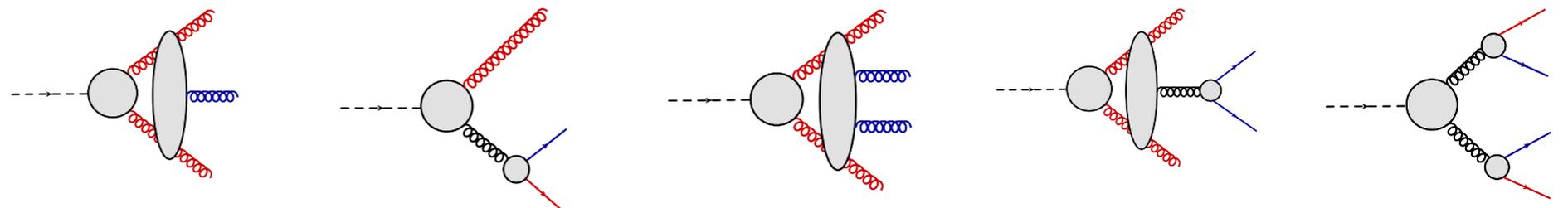
Quark-gluon:

$$\bar{\chi}^0 \rightarrow \tilde{g}g + \dots$$



Gluon-gluon:

$$H \rightarrow gg + \dots$$



# Where to start?

$pp \Rightarrow 2j$  at NNLO is a complicated calculation:

- many crossings and colour factors to consider
- up to four massless partons in the final state means a large number of (overlapping) unresolved limits

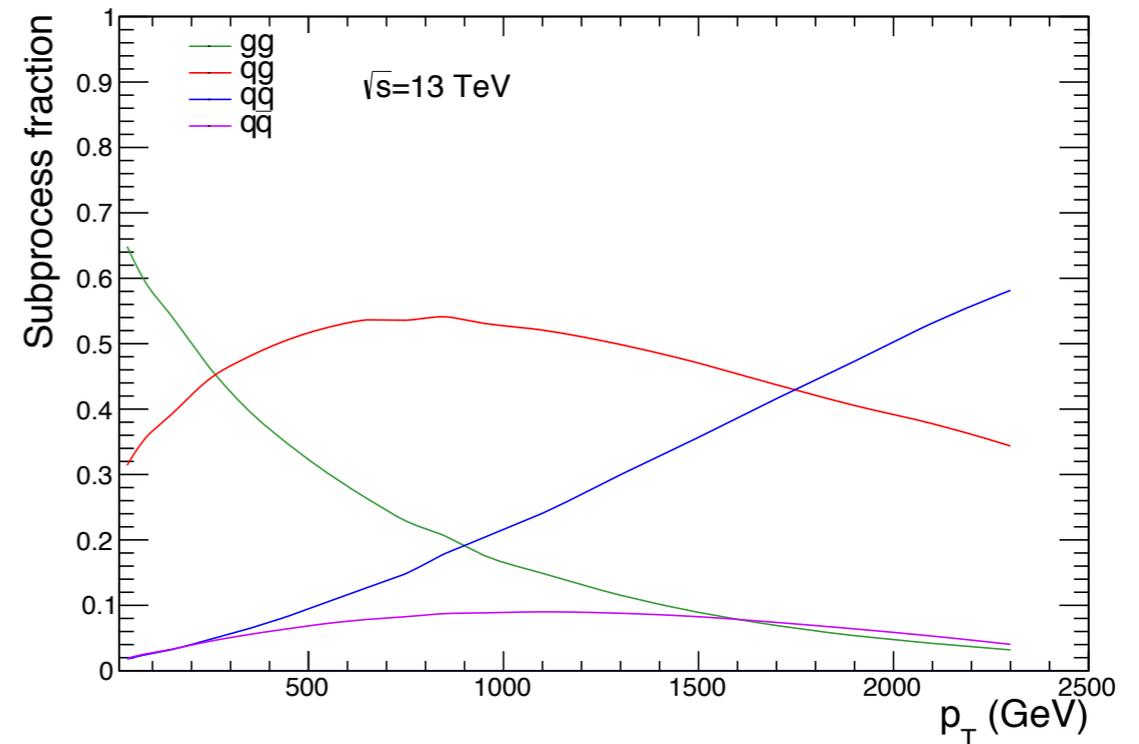
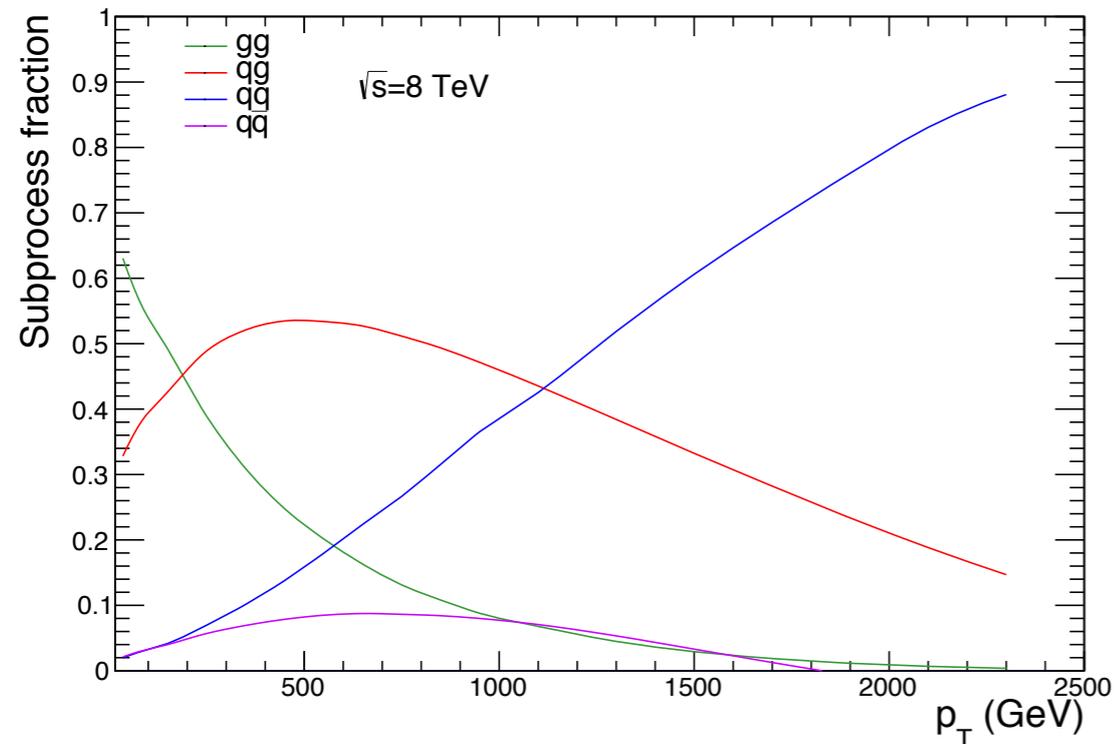
Start by considering:

- what are the most important channels?
- what are the most important colour factors in each channel?

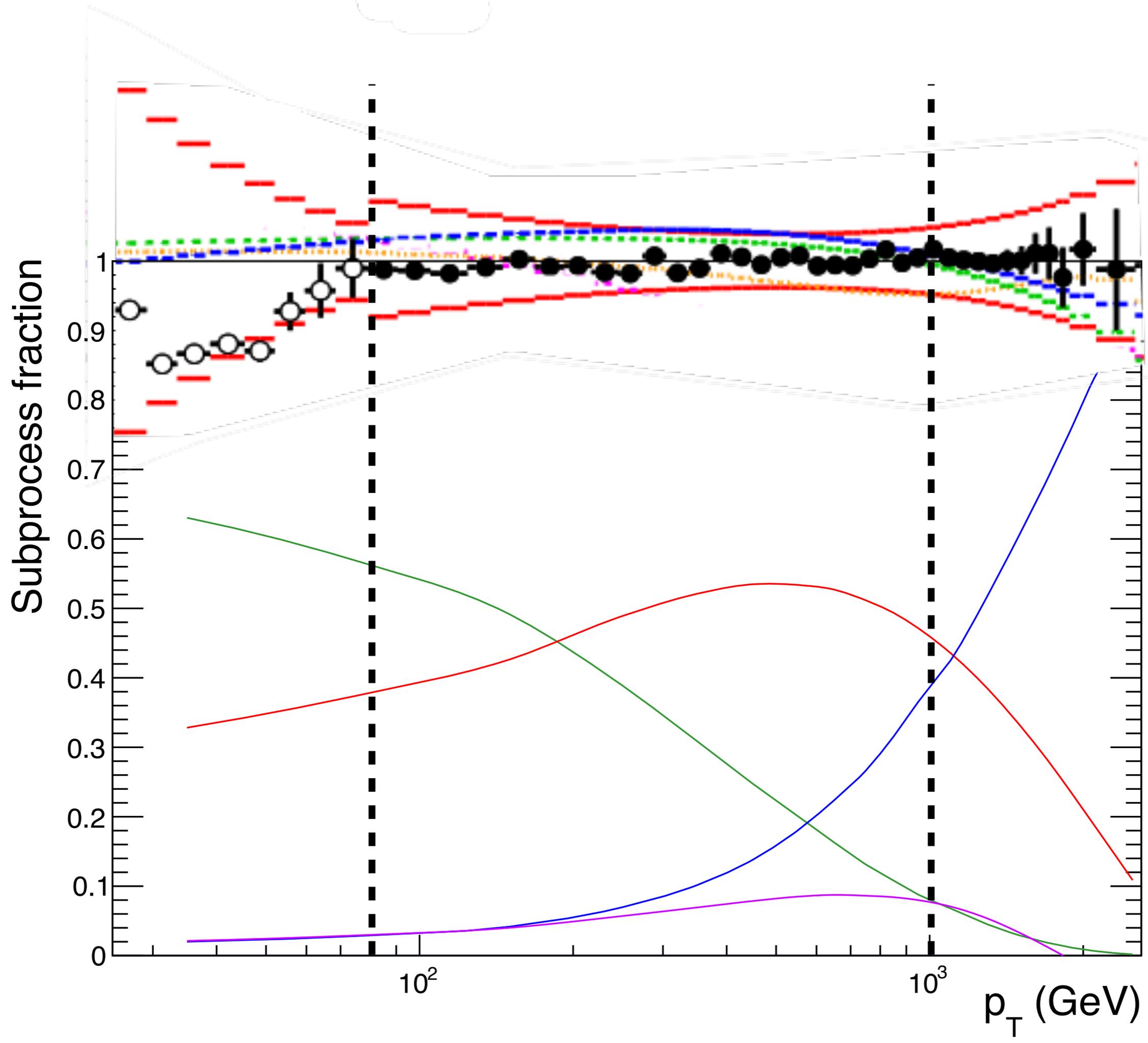
# Channels

At low to moderate  $p_T$  the gluonic initial-states ( $gg+qg$ ) dominate

At high  $p_T$  quark scattering becomes important



In this talk we will focus on  $gg+qg$ ;  $qq$  results in preparation



# gg channel

Start with the double real all-gluon contribution [Glover, Pires '10]:

- six gluon matrix element [Mangano, Parke, Xu '87; Berends, Giele '87]:

$$|\mathcal{M}_6^0|^2 = \sum_{\text{perms}} A_6^0(1, 2, i, j, k, l)$$

- “single unresolved” subtraction term:

$$f_3^0(2, i, j) A_5^0(1, \bar{2}, (\widetilde{ij}), k, l)$$

- “double unresolved” subtraction term:

$$F_4^0(2, i, j, k) A_4^0(1, \bar{2}, (\widetilde{ijk}), l)$$

- “spurious unresolved” subtraction term:

$$f_3^0(2, i, j) f_3^0((\widetilde{ij}), k, l) A_4^0(1, \bar{2}, (\widetilde{(\widetilde{ij})(kl))})$$

Real-virtual correction [Glover, Pires '12]:

- one-loop five gluon matrix element [Bern, Dixon, Kosower '93]

$$|\mathcal{M}_5^1|^2 = \sum_{\text{perms}} A_5^1(1, 2, i, j, k)$$

- pole subtraction term:

$$\mathbf{J}^{(1)}(1, 2, i, j, k) A_5^0(1, 2, i, j, k)$$

- single unresolved subtraction term:

$$f_3^0(2, i, j) A_4^1(1, \bar{2}, (\tilde{i}j), k) + f_3^1(2, i, j) A_4^0(1, \bar{2}, (\tilde{i}j), k)$$

- spurious pole/single unresolved subtraction term:

$$\mathbf{J}^{(1)}(2, i) f_3^0(i, j, k) A_4^0(1, \bar{2}, (\tilde{i}j), (\tilde{j}k))$$

## Double virtual correction [Gehrmann, Gehrmann de-Ridder, Glover, Pires '13]

- two-loop matrix elements [Glover, Oleari, Tejeda-Yeomans '01]

• 1:

$$|\mathcal{M}_4^2|^2 = \sum_{\text{perms}} A_4^2(1, 2, i, j)$$

- double virtual subtraction term:

$$\begin{aligned} &+ \mathbf{J}^{(1)}(1, 2, i, j) A_4^1(1, 2, i, j) \\ &+ \frac{1}{2} \mathbf{J}^{(1)}(1, 2, i, j) \otimes \mathbf{J}^{(1)}(1, 2, i, j) A_4^0(1, 2, i, j) \\ &+ \mathbf{J}^{(2)}(1, 2, i, j) A_4^0(1, 2, i, j) \end{aligned}$$

- analogous to well known IR pole structure [Catani '98]
- structure is universal and generalizable to higher multiplicities

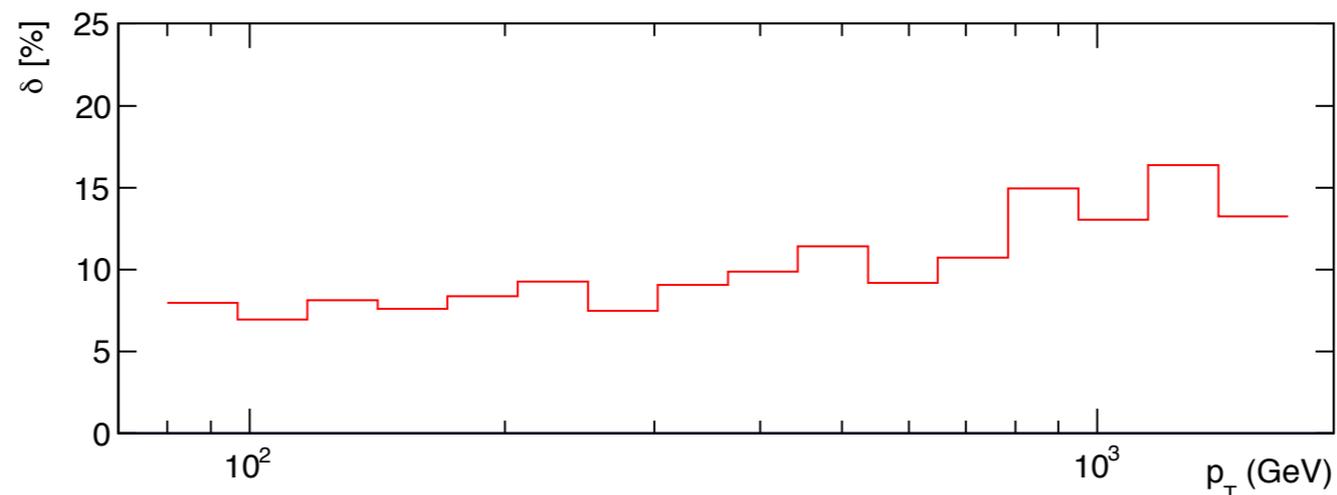
Sub-leading colour all-gluon correction [JC, Glover, Pires '14]:

- posed an interesting theoretical challenge
- antenna subtraction designed for squared partial amplitudes
- sub-leading colour RR and RV gluon scattering built from interferences

$$\mathcal{A}_6^{0,\dagger}(\sigma) [\mathcal{A}_6^0(\sigma') + \mathcal{A}_6^0(\sigma'') + \mathcal{A}_6^0(\sigma''')]$$

$$\mathcal{A}_5^{0,\dagger}(\sigma) \mathcal{A}_5^1(\sigma')$$

The method worked well and produced a small correction



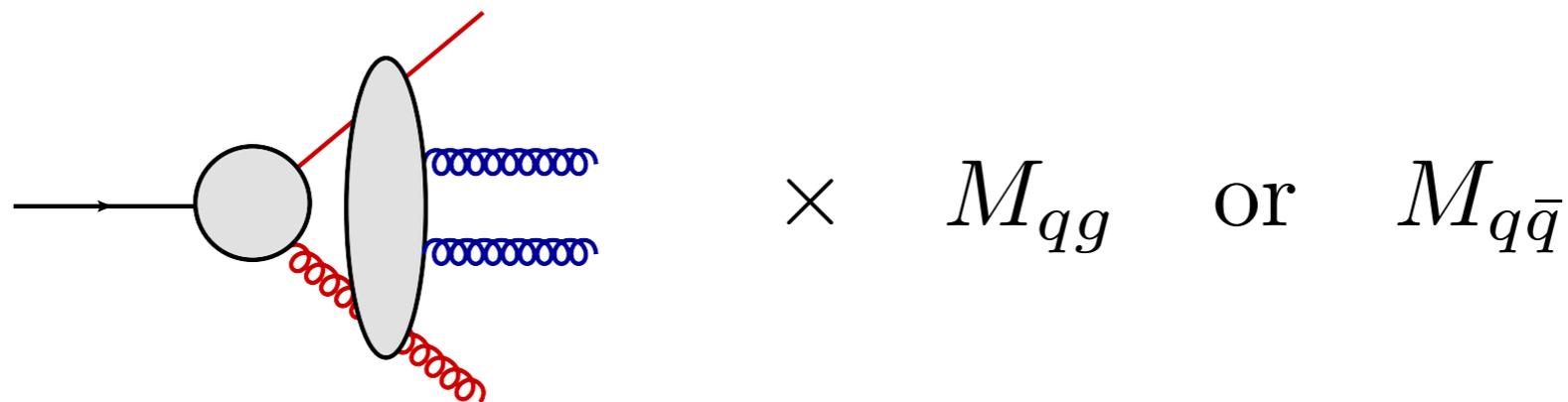
# qg channel

Very important channel over a wide range of  $p_T$ :

- main missing component for jets up to  $\sim 1$  TeV

Also presents interesting theoretical challenges:

- antennae interpolate between many limits with a smooth momentum map
- not always desirable when factoring onto physically different matrix elements



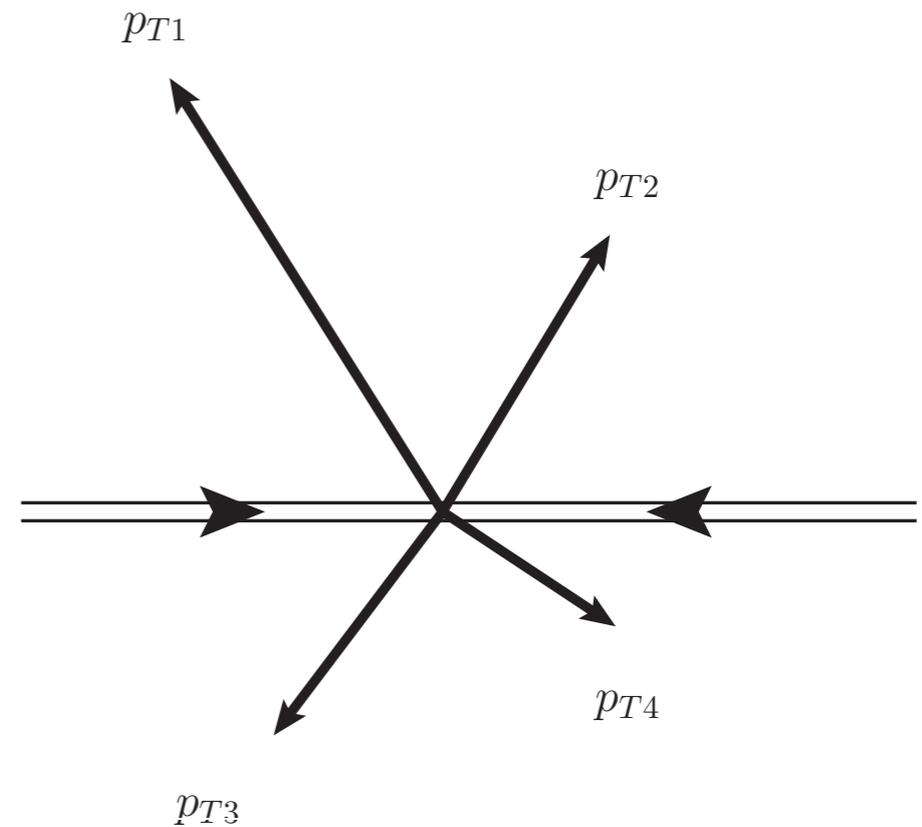
Limits can be disentangled successfully and systematically

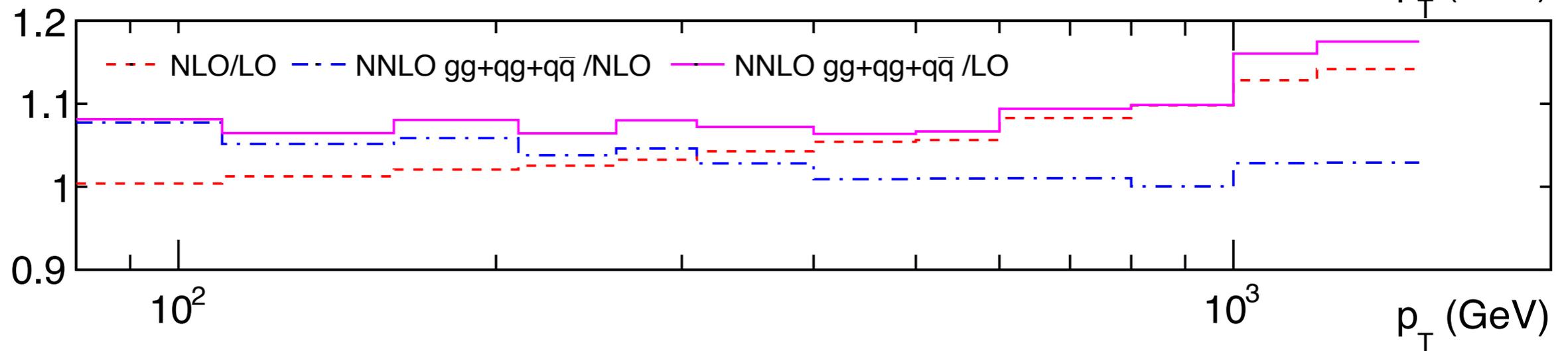
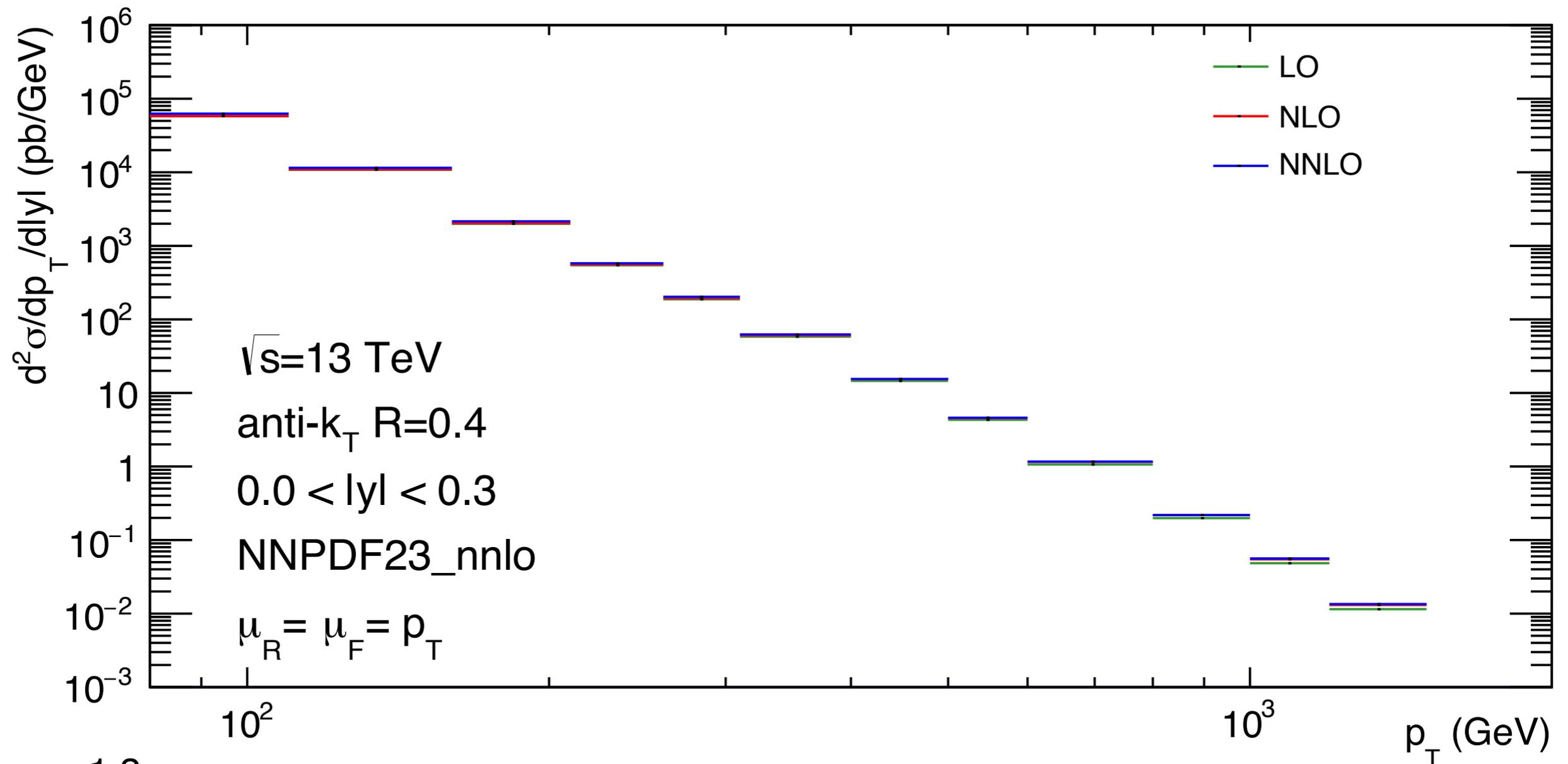
# Results

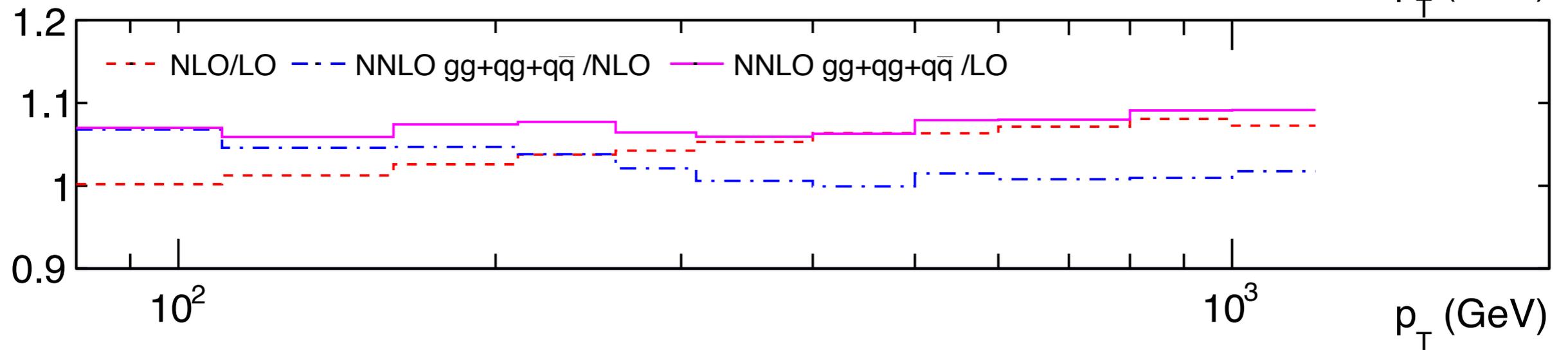
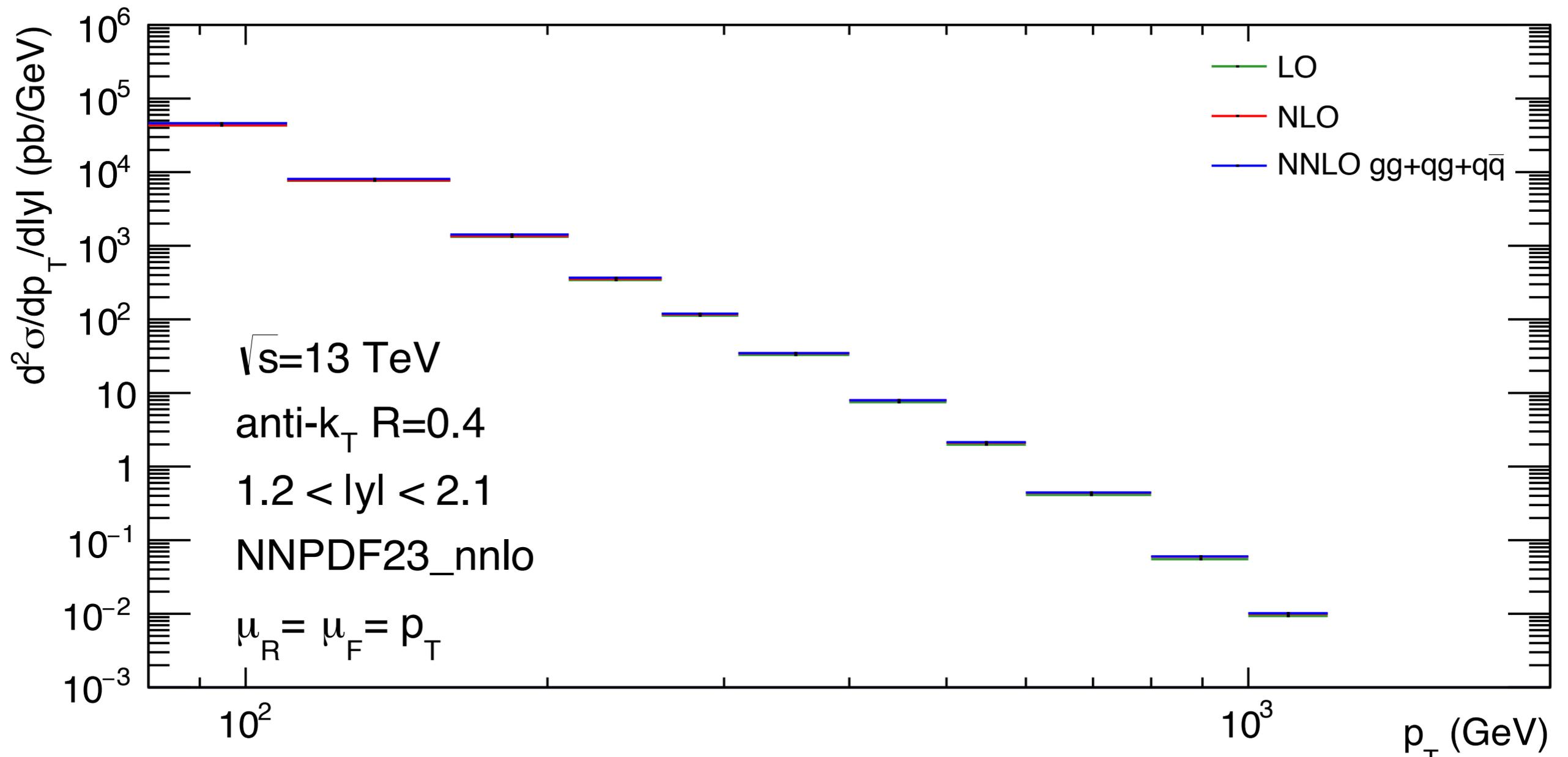
The following results are for  $gg+qg+q\bar{q}\Rightarrow 2j$  at 13 TeV

Setup:

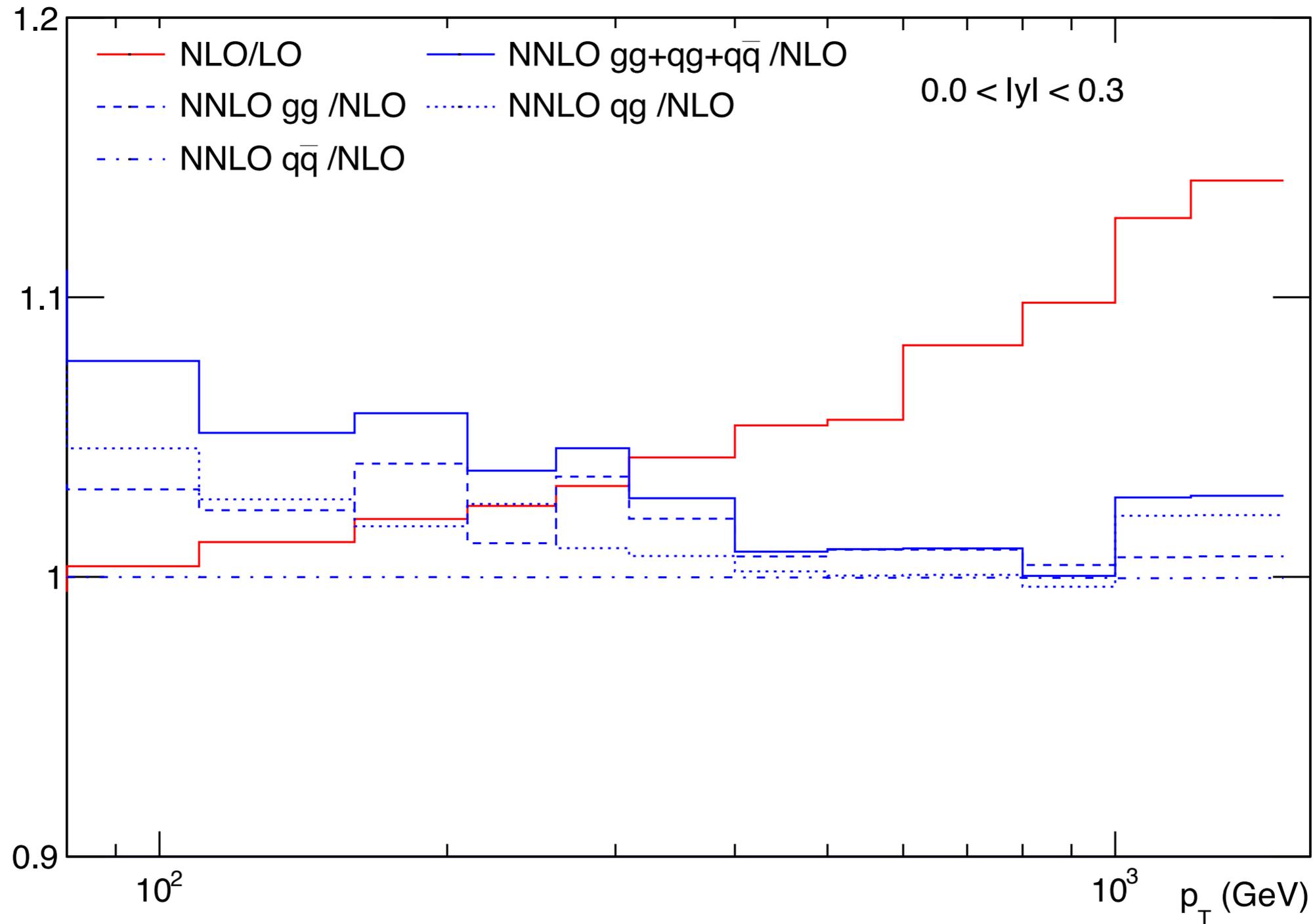
- NNPDF2.3\_NNLO
- accept jets with  $p_T > 80$  GeV
- rapidity cut  $|y| < 4.4$
- scale  $\mu = \mu_F = p_T$  rather than  $p_{T1}$
- anti- $k_T$  jet algorithm  $R=0.4$

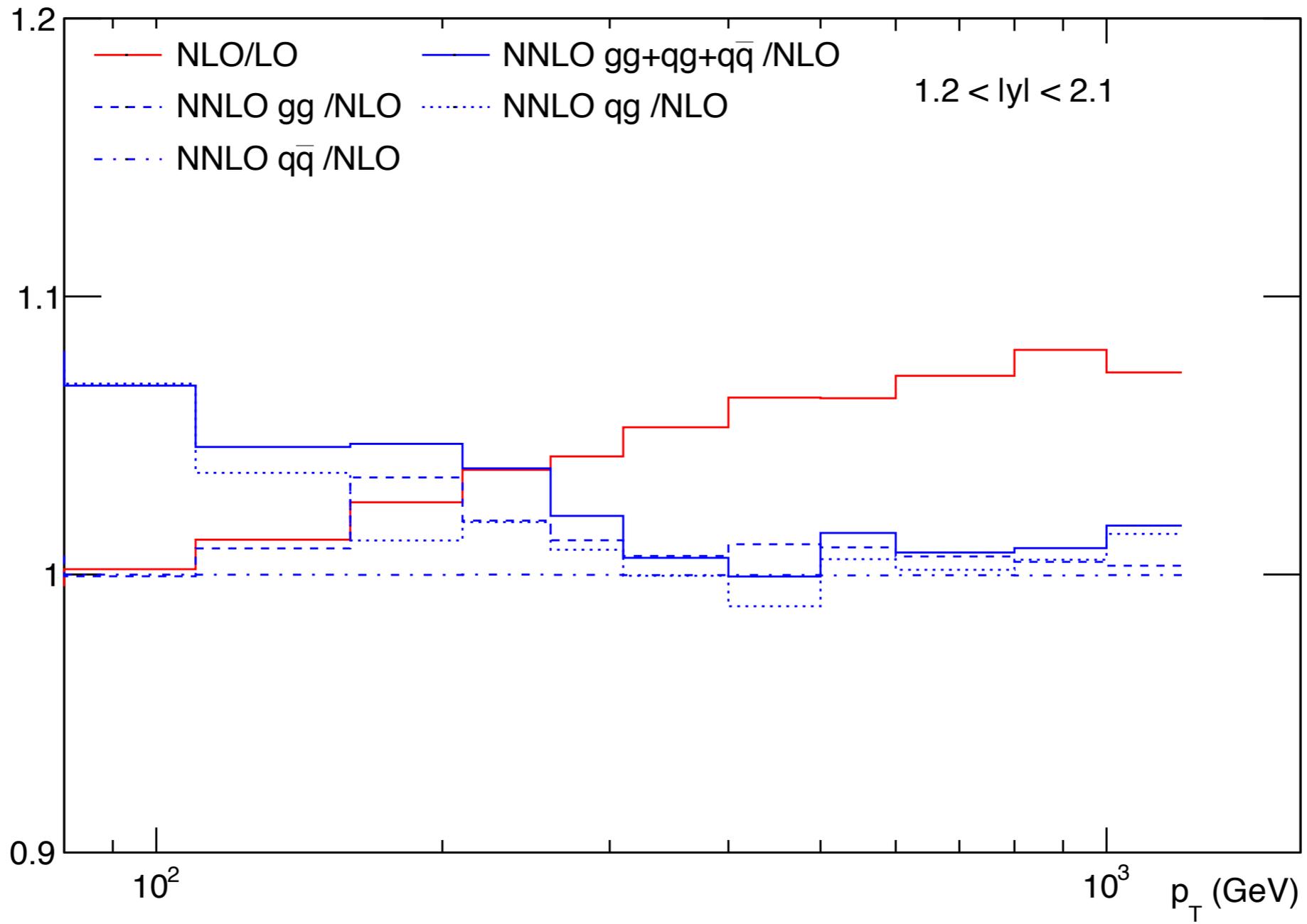


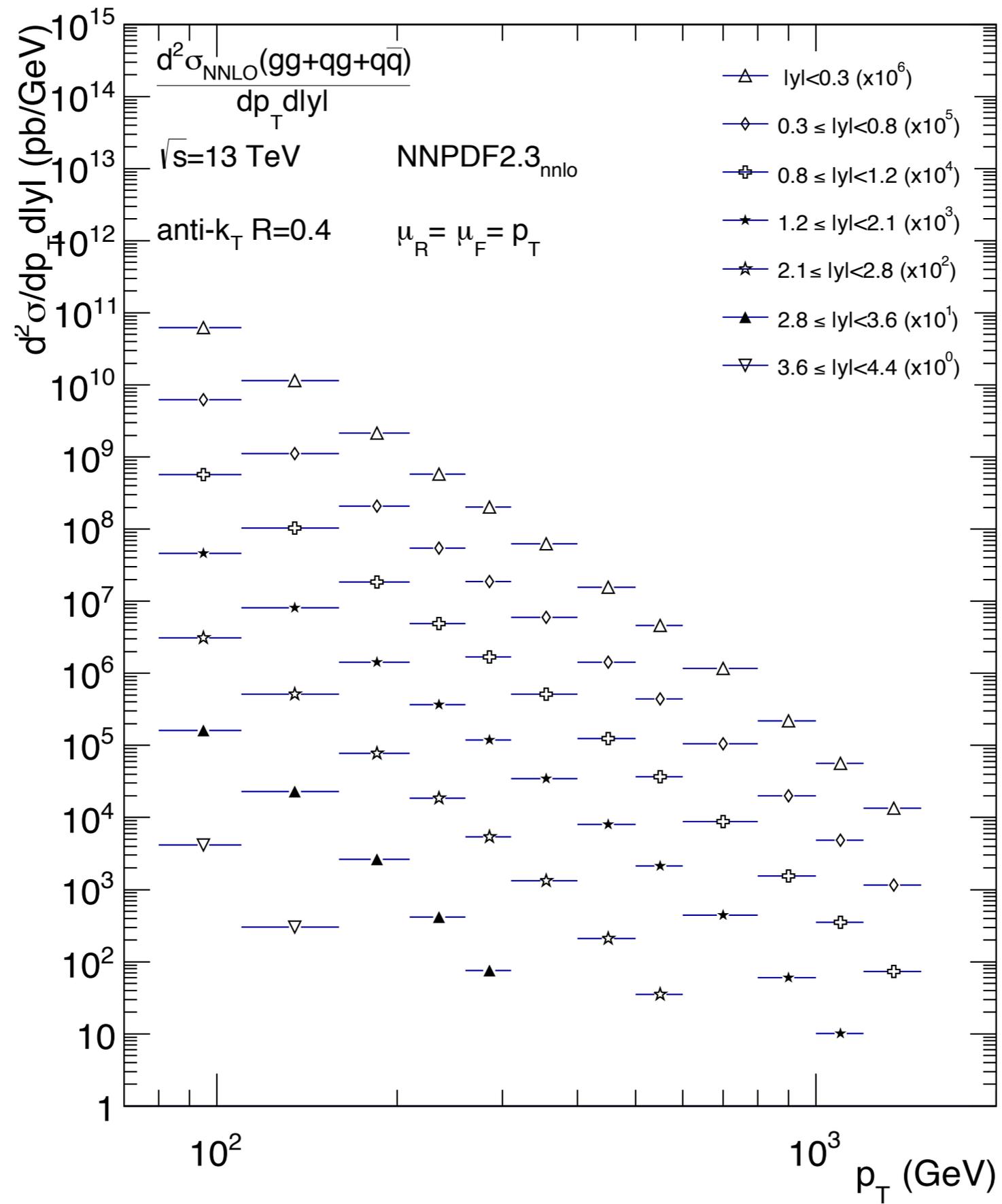




# K-factors







# Where to go now?

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 5 people like this.



**Joey Huston** 2066. So that's the final prediction when the inclusive jet cross section will be finished?

Like · Reply ·  2 · 11 June at 08:36



A few things remain to be included in our study:

- quark scattering for high  $p_T$  jets (results in preparation)
- leading  $N_F$  corrections in all channels (in preparation)
- sub-leading colour probably insignificant and can be dropped without compromising phenomenology
- updated scale variation (in preparation)
- upgrade of the Monte Carlo and interface to APPLgrid/n-tuples

# Summary

Antenna subtraction is a flexible and powerful IR subtraction scheme

We have used this method to calculate the NNLO correction to jet production at the LHC:

- updated gluon scattering results at 13 TeV
- added the new and significant quark-gluon channel
- we observe that the new qg channel dominates for moderate  $p_T$
- NNLO corrections up to  $\sim 8\%$  of the NLO decreasing with  $p_T$
- quark-quark channel and  $N_F$  results in preparation