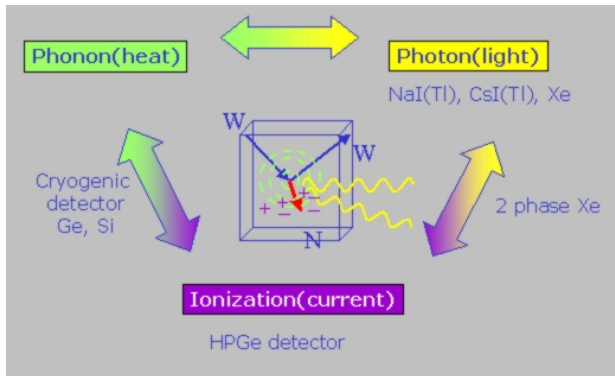


Beyond Annual Modulation at WIMP Direct-Detection Experiments

Samuel Lee
Princeton Center for Theoretical Science
PACIFIC 2014
September 15, 2014

- Basics of WIMP direct detection & annual modulation
- CoGeNT and DAMA modulation signals
- A more detailed treatment of modulation?
 - Gravitational focusing of WIMPs (and cosmic neutrinos)
 - Higher-harmonic modes

Direct Detection 101



Direct detection of **WIMP-induced nuclear recoils** via phonons, scintillation, ionization, etc.

Event rate given by

$$\frac{dR}{dE}(t) = \underbrace{\frac{\sigma_0 F^2(E)}{2m_\chi \mu_{\chi N}^2}}_{\text{PARTICLE}} \times \rho_0 \underbrace{\int_{v_{\min}}^{\infty} d^3v \frac{f(\mathbf{v}, t)}{v}}_{\text{ASTRO}}$$

Event rate given by

$$\begin{aligned}\frac{dR}{dE}(t) &= \underbrace{\frac{\sigma_0 F^2(E)}{2m_\chi \mu_{\chi N}^2}}_{\text{PARTICLE}} \times \rho_0 \underbrace{\int_{v_{\min}}^{\infty} d^3v \frac{f(\mathbf{v}, t)}{v}}_{\text{ASTRO}} \\ &= \frac{\sigma_0 F^2(E)}{2m_\chi \mu_{\chi N}^2} \times \rho_0 \underbrace{\eta(v_{\min}, t)}_{\text{mean inverse speed}}\end{aligned}$$

Event rate given by

$$\begin{aligned}
 \frac{dR}{dE}(t) &= \underbrace{\frac{\sigma_0 F^2(E)}{2m_\chi \mu_{\chi N}^2}}_{\text{PARTICLE}} \times \underbrace{\rho_0 \int_{v_{\min}}^{\infty} d^3v \frac{f(\mathbf{v}, t)}{v}}_{\text{ASTRO}} \\
 &= \frac{\sigma_0 F^2(E)}{2m_\chi \mu_{\chi N}^2} \times \rho_0 \underbrace{\eta(v_{\min}, t)}_{\text{mean inverse speed}}
 \end{aligned}$$

Minimum lab-frame WIMP speed required to induce nuclear recoil with energy E

$$v_{\min} = \sqrt{\frac{m_N E}{2\mu_{\chi N}^2}}$$

Event rate depends on lab-frame WIMP velocity distribution

$$f(\mathbf{v}, t) = g(\mathbf{v} + \mathbf{v}_{\text{lab}}(t))$$

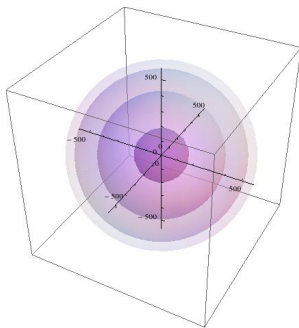
Event rate depends on lab-frame WIMP velocity distribution

$$f(\mathbf{v}, t) = g(\mathbf{v} + \mathbf{v}_{\text{lab}}(t))$$

Standard Halo Model (truncated Maxwellian) generally assumed for galactic-frame velocity distribution

$$g(\mathbf{v}) \propto e^{-v^2/v_0^2} \theta(v_{\text{esc}} - v)$$

$$v_0 = 220 \text{ km/s}, \quad v_{\text{esc}} \approx 550 \text{ km/s}$$



Event rate depends on lab-frame WIMP velocity distribution

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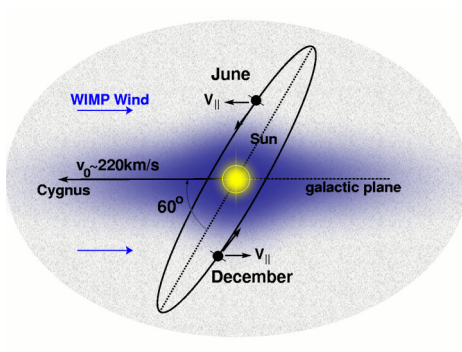
Event rate can be written as

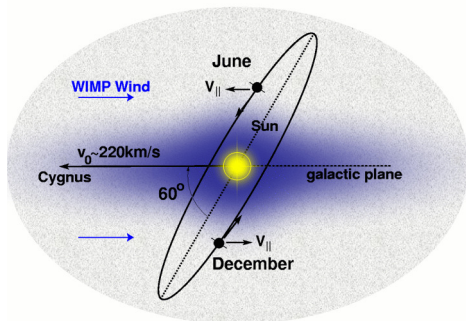
$$\frac{dR}{dE} \propto \int_{v_{\text{min}}(E)}^{\infty} d^3v \frac{f(\mathbf{v})}{v} \approx A_0(E)$$

For SHM, and typical WIMP and experiments,

$$A_0(E) \propto e^{-E/E_0}, \quad E_0 \approx 10 \text{ keV}$$

Modulation 101





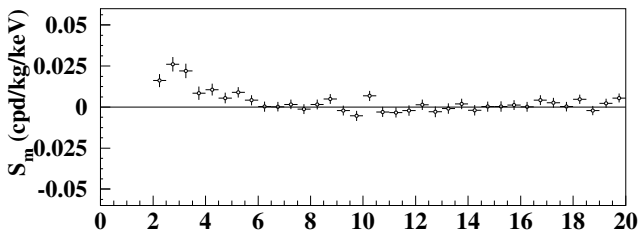
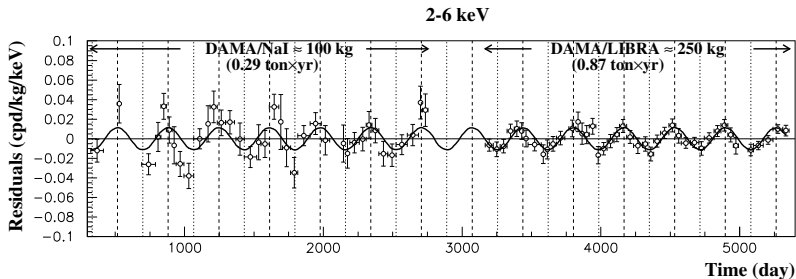
Modulation of event rate can be written as

$$\frac{dR}{dE}(t) \propto \int_{v_{\min}(E)}^{\infty} d^3v \frac{f(\mathbf{v}, t)}{v} \approx A_0(E) + A_1(E) \cos[\omega(t - t_{\max})]$$

For SHM, and typical WIMP and experiments,

$$A_0(E) \propto e^{-E/E_0}, \quad E_0 \approx 10 \text{ keV}, \quad \bar{A}_1/\bar{A}_0 \approx \text{few}\%, \quad t_{\max} \approx \text{June 1}$$

Modulation at DAMA: Signal

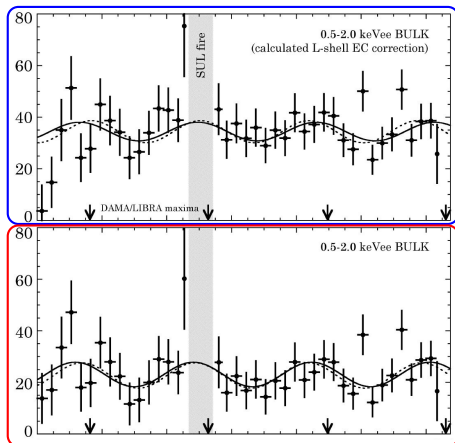


$$T = 0.999 \pm 0.002 \text{ year}, \quad t_{\max} = \text{May } 25 \pm 7 \text{ days}$$

Modulation at DAMA: Summary

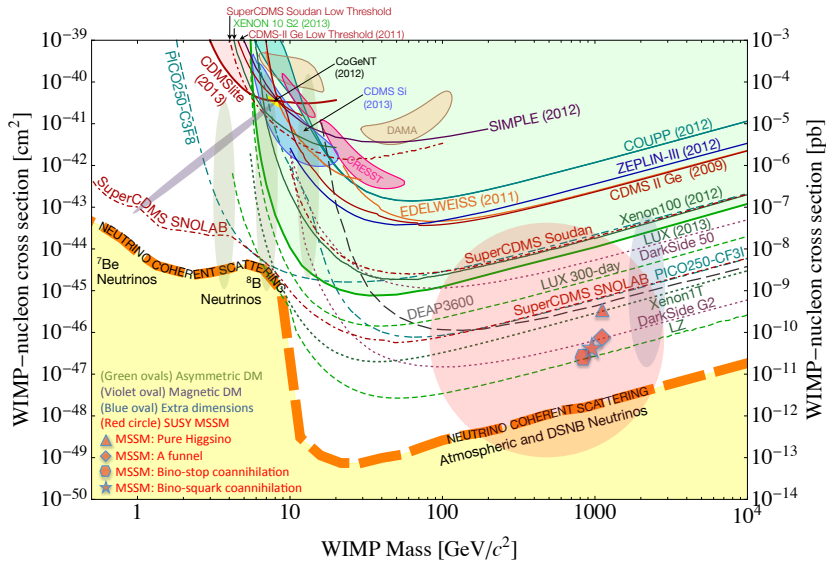
- Exposure of \sim ton-year at NaI target
- $\sim 9\sigma$ annual modulation of single-hit residuals at 2–6 keVee
- Assuming SHM, 10 GeV and 80 GeV ROIs allowed
- Phase compatible, fractional amplitude not measured

- Calculating EC BG:
 - $T = 336 \pm 24$ days
 - $S = 12.4 \pm 5\%$
 $\rightarrow \bar{A}_1/\bar{A}_0 \approx 35\%$
- Fitting EC BG:
 - $T = 350 \pm 20$ days
 - $S = 21.7 \pm 15\%$
 $\rightarrow \bar{A}_1/\bar{A}_0 \approx 62\%$
- Fixing $T = 1$ year:
 $t_{\max} = \text{April } 13 \pm 47$ days



- Exposure of 1,129 live days at Ge target
- Events in low-energy (0.5–2.0 keVee) and high-energy (2.0–4.5 keVee) bins
- Improved bulk/surface discrimination via rise-time cuts
- $\sim 2.2\sigma$ annual modulation in low-energy, bulk events (decrease from $\sim 2.8\sigma$)
- For SHM, best fit by ~ 8 GeV & 2×10^{-41} cm²
- Phase compatible, but fractional amplitude ~ 4 -7x larger

Exclusion at Other Experiments?



Billard et al. 2013 (and others...)

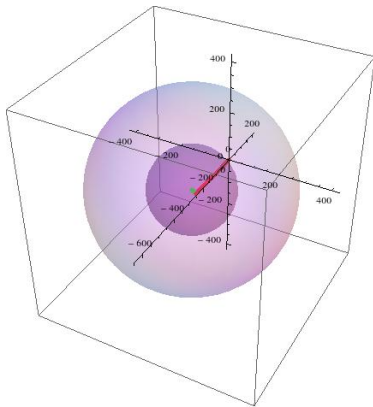
- Making modulation a better smoking gun?
- Moving towards a holistic understanding of dark-matter particle physics and astrophysics
- Sensitivity to velocity substructures (disk, streams, etc.)
- May extend reach below neutrino floor (if necessary...)

Galilean transformation relates lab and galactic frame

$$\begin{aligned} f(\mathbf{v}, t) &= g(\mathbf{v} + \mathbf{v}_{\text{lab}}(t)) \\ &= g(\mathbf{v} + \mathbf{v}_{\odot} + \mathbf{v}_{\oplus}(t)) \end{aligned}$$

Event rate $\frac{dR}{dE}(t)$ proportional to integral of

$$\frac{f(\mathbf{v}, t)}{v}$$



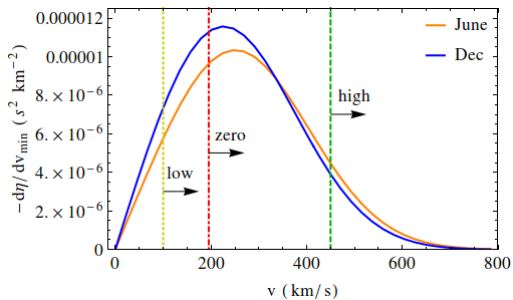
Event rate $\frac{dR}{dE}(t)$ proportional to integral of

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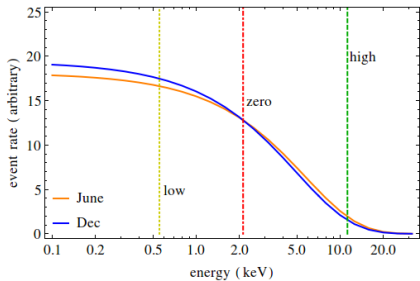
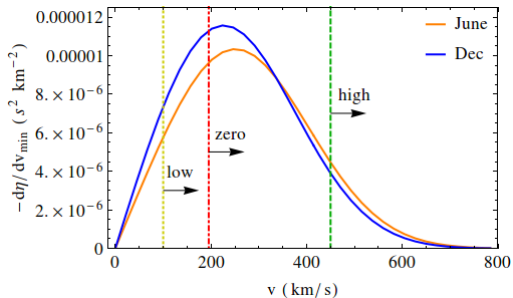
Event rate $\frac{dR}{dE}(t)$ proportional to integral of

$$\frac{f(\mathbf{v}, t)}{v}$$

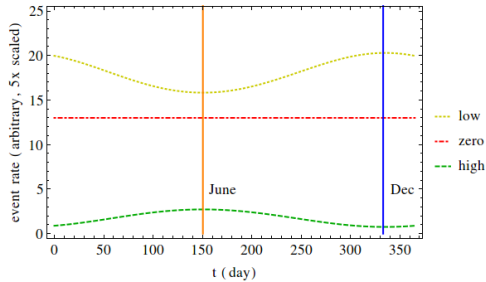
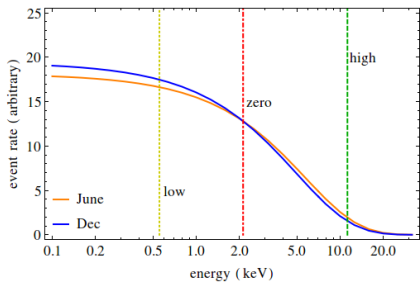
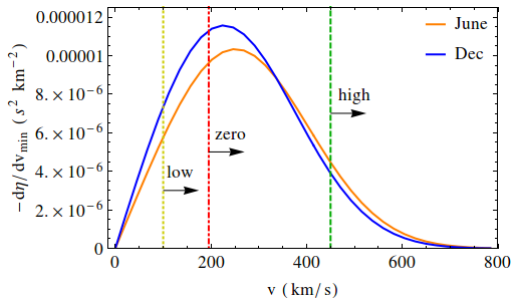
integrated over $v \geq v_{\min}(E)$



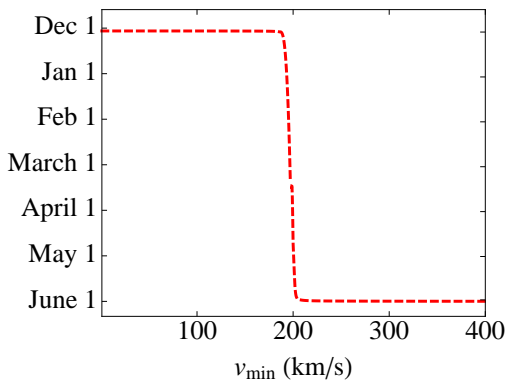
Modulation 102



Modulation 102



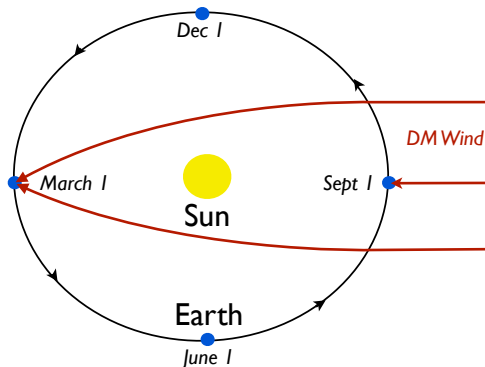
Phase of the modulation flips: maximum of event rate in December (June) at low (high) v_{\min}



Modulation 201: Gravitational Focusing

Sun's gravitational potential induces comparable modulation
(for slowly moving WIMPs)

$$\bar{A}_1/\bar{A}_0 \sim \epsilon \sim \frac{v_{\oplus}}{v_{\odot}} \sim \frac{\sqrt{GM_{\odot}/R_{\oplus}}}{v_0}$$

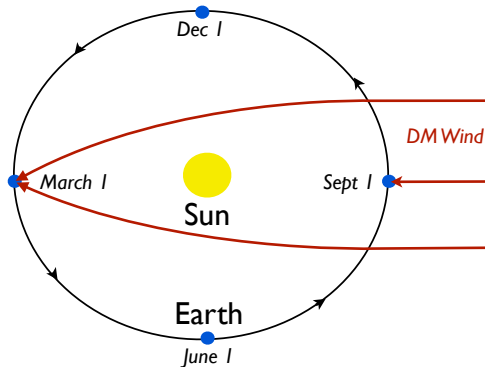


Modulation 201: Gravitational Focusing

joepoppa

January 6, 2014, 2:35 PM

At long last, we know something definite about dark matter: it hibernates, during the winter. That, alone, should yield a wealth of clues. I think the reason dark matter is most noticeable, come March, is that long hibernation has made it VERY hungry. What they should now do is launch a thorough investigation into bear hibernation, and that might afford more clues into dark matter hibernation.

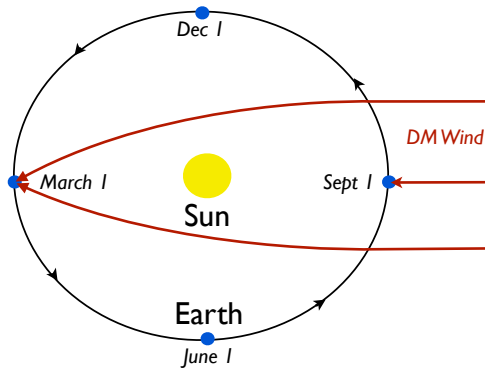


SKL, Mariangela Lisanti, Annika Peter, Ben Safdi 2013

Modulation 201: Gravitational Focusing

Non-Galilean transformation relates lab and galactic frames

$$f(\mathbf{v}, t) = g(\mathbf{v}_\infty [\mathbf{v} + \mathbf{v}_\oplus(t), t] + \mathbf{v}_\odot)$$
$$\mathbf{v}_\infty[\mathbf{v}_s, t] = \frac{v_\infty^2 \mathbf{v}_s + v_\infty (G M_\odot / r_s) \hat{\mathbf{r}}_s - v_\infty \mathbf{v}_s (\mathbf{v}_s \cdot \hat{\mathbf{r}}_s)}{v_\infty^2 + G M_\odot / r_s - v_\infty (\mathbf{v}_s \cdot \hat{\mathbf{r}}_s)}$$



Modulation 201: Gravitational Focusing

Non-Galilean transformation relates lab and galactic frames

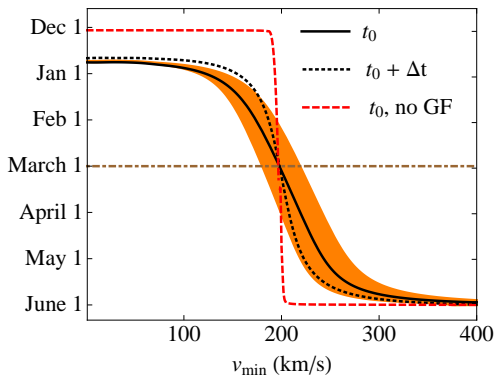
$$\begin{aligned} f(\mathbf{v}, t) &= g(\mathbf{v}_\infty [\mathbf{v} + \mathbf{v}_\oplus(t), t] + \mathbf{v}_\odot) \\ \mathbf{v}_\infty[\mathbf{v}_s, t] &= \frac{v_\infty^2 \mathbf{v}_s + v_\infty (G M_\odot / r_s) \hat{\mathbf{r}}_s - v_\infty \mathbf{v}_s (\mathbf{v}_s \cdot \hat{\mathbf{r}}_s)}{v_\infty^2 + G M_\odot / r_s - v_\infty (\mathbf{v}_s \cdot \hat{\mathbf{r}}_s)} \end{aligned}$$

Modulation 201: Gravitational Focusing

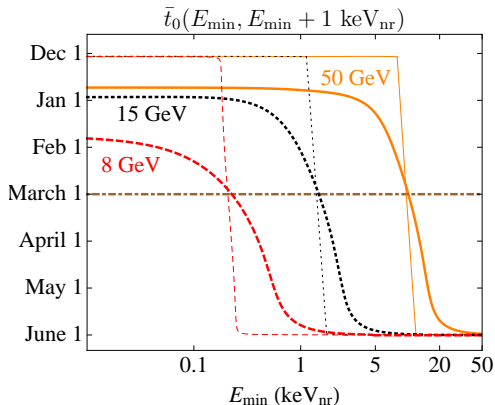
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Gravitational focusing results in an
ENERGY-DEPENDENT MODULATION PHASE

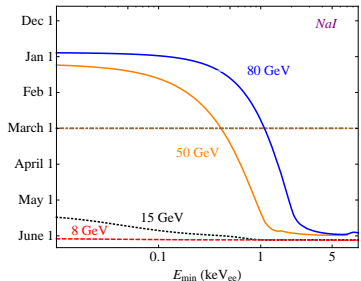
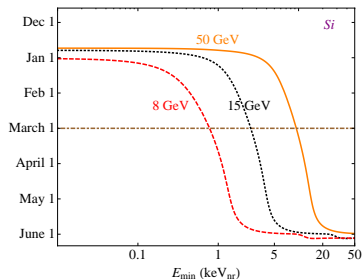
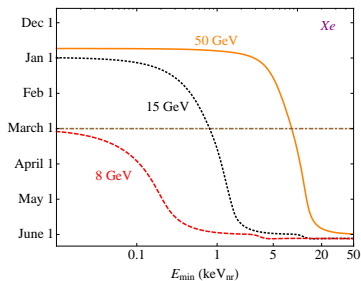
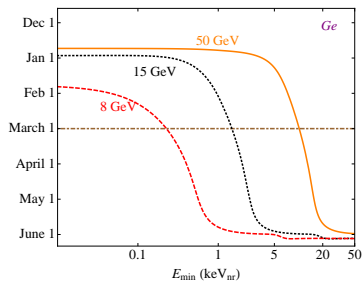


Gravitational focusing results in an
ENERGY-DEPENDENT MODULATION PHASE

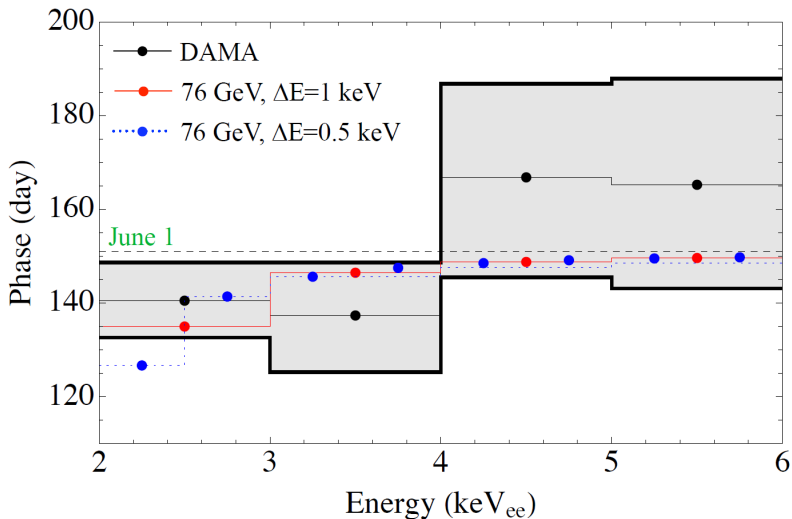


Smoking gun 2.0 – but **low energy thresholds** required!

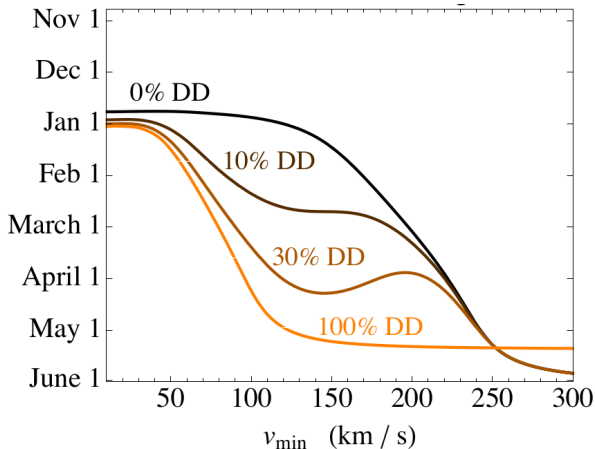
Modulation 201: Gravitational Focusing



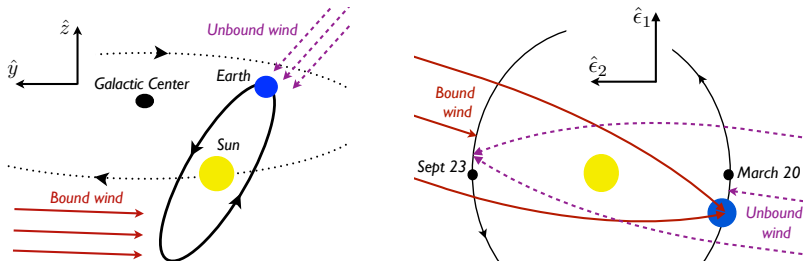
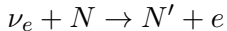
GF may already be important for 80 GeV DAMA ROI



Co-rotating dark-disk scenarios with slowly moving WIMPs will be especially sensitive to GF



GF is **solely responsible** for modulation of cosmic neutrinos, which may be detected via **velocity-independent** capture



Safdi, Lisanti, Spitz, Formaggio 2014

Modulation 201: Higher Harmonics

Higher modes are generically present, but suppressed by $\epsilon = \frac{v_{\oplus}}{4v_{\odot}}$
Circular orbit usually assumed, leading to cosine expansion

$$\frac{dR}{dE}(t) = A_0(E) + \sum_n A_n(E) \cos[n\omega(t - t_{\max})]$$

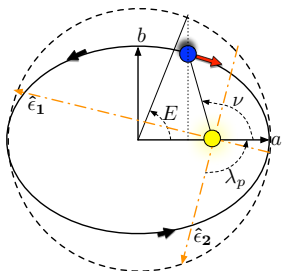
with $A_n/A_0 \sim \epsilon^n$ and $MT(A_n)/MT(A_0) \sim \epsilon^{-2n}$

Modulation 201: Higher Harmonics

Consistent expansion should include **eccentricity** $e = 0.017 \sim \epsilon$
Elliptical orbit (correcting Lewin and Smith 1996) leads to

$$\frac{dR}{dE}(t) = A_0(E) + \sum_n A_n(E) \cos[n\omega(t - t_{\max})] \\ + \sum_n B_n(E) \sin[n\omega(t - t_{\max})]$$

with $B_n/A_0 \sim \epsilon^n$ and $MT(B_n)/MT(B_0) \sim \epsilon^{-2n}$ (except $n = 1$)

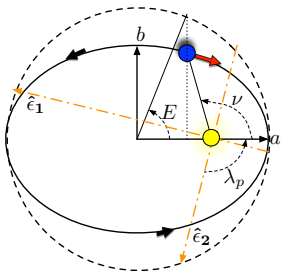


Modulation 201: Higher Harmonics

Harmonic structure is typically related to orbital parameters

$$\frac{B_1(E)}{A_1(E)} \approx 2e \sin(\lambda_p - w\phi) \approx \frac{1}{59}$$

$$\frac{B_2(E)}{B_1(E)} \approx \frac{1}{2}$$



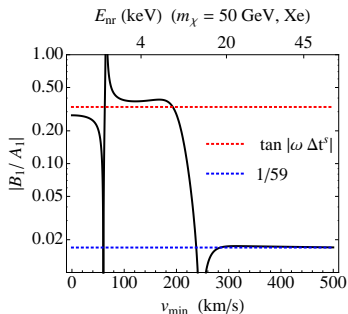
Modulation 201: Higher Harmonics

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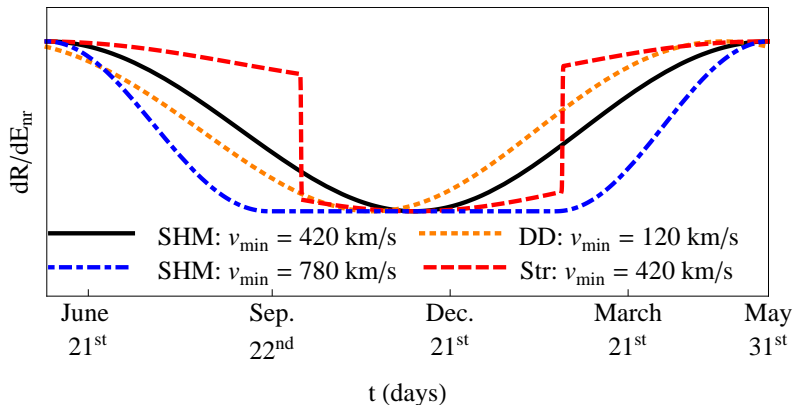
$$\frac{B_2(E)}{B_1(E)} \approx \frac{1}{2}$$

Relations can be used to test presence of velocity substructure

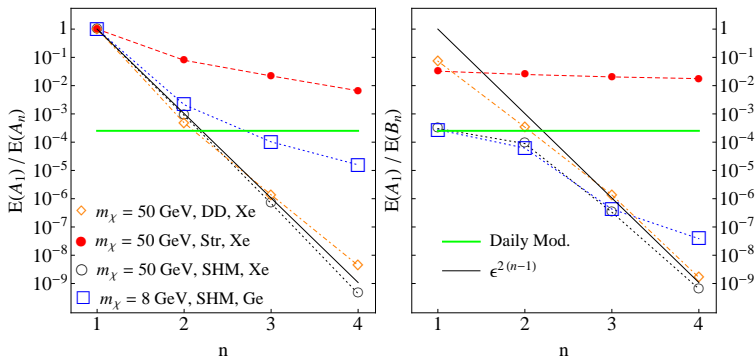


Modulation 201: Higher Harmonics

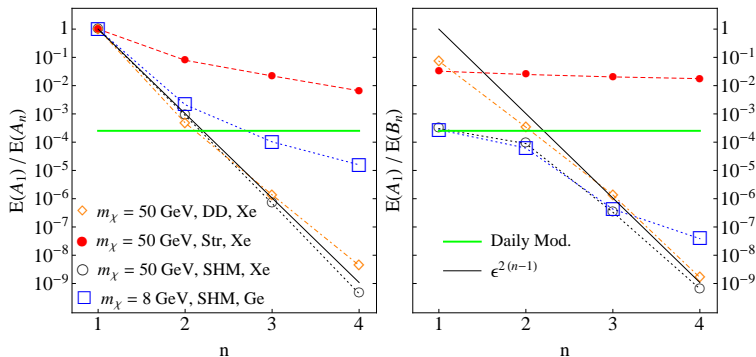
Higher harmonics can be enhanced in various scenarios



Higher harmonics can be enhanced in various scenarios



Higher harmonics can be enhanced in various scenarios



Orbit of the earth and DM velocity distribution both leave
UNIQUE IMPRINT ON HARMONIC STRUCTURE

- Confusion surrounding hints of signals motivates a more detailed consideration of modulation
- Higher-order effects can be powerful and informative discriminators
 - Gravitational focusing by sun → energy-dependent phase
 - Details of earth's orbit → structure of higher harmonics
- Low thresholds (GF) or large exposures (harmonics)
- Time dependence provides a key axis (along with directional and material signals) for DM detection and characterization!