

A WARPED MODEL OF DARK MATTER

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Australia

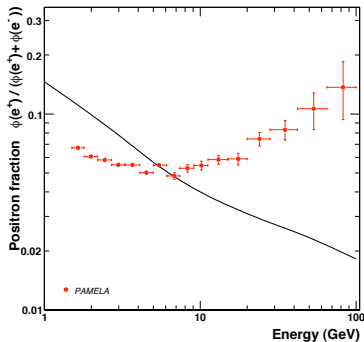
PACIFIC 2011

Based on work
with Tony Gherghetta

arXiv:1002.2967

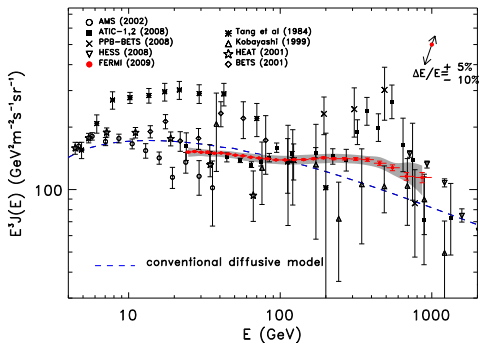
INTRODUCTION: COSMIC RAY ANOMALIES

- PAMELA found rising positron fraction $\frac{\phi(e^+)}{\phi(e^+) + \phi(e^-)}$ in energy range 10-100 GeV
- FERMI and HESS saw excess (over expected background) in $e^+ + e^-$ flux in range 100-~1000 GeV



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Are these signals due to **annihilating dark matter**?

Challenge for model builders:

- PAMELA saw no excess in antiproton flux
- Annihilation cross section to explain cosmic ray signals is $\mathcal{O}(100 - 1000)$ larger than cross section for right relic abundance

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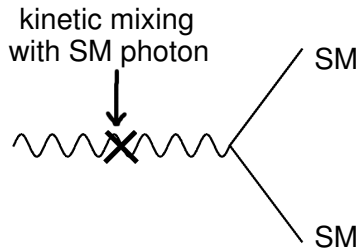
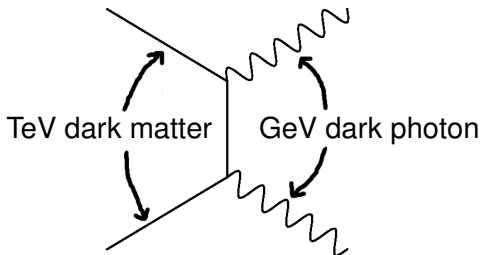
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OUTLINE OF THE TALK

1 GEV GAUGE BOSONS FROM WARPED SPACE

2 MIXING WITH STANDARD MODEL PHOTON

3 DARK MATTER

- Dark matter on the IR brane
- Dark Matter in the bulk
- Gauge couplings

4 CONCLUSIONS

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Where does the **GeV scale** come from?

- SUSY (KATZ & SUNDRUM; MORISSEY, POLAND & ZUREK; ...)
- We use: **warped extra dimension** (see also McDONALD & MORISSEY)

Idea:

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The diagram shows two vertical lines representing branes. The left line is labeled "Planck brane" and the right line is labeled "TeV brane". A horizontal line connects the two vertical lines, representing the extra dimension. In the center of this region, the metric is given by the equation:

$$ds^2 = e^{-2ky} (\eta_{\mu\nu} dx^\mu dx^\nu) + dy^2$$

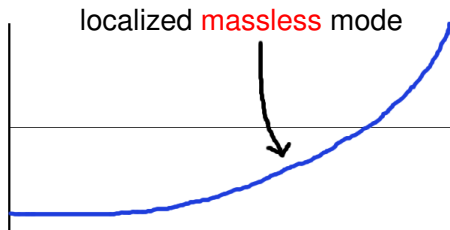
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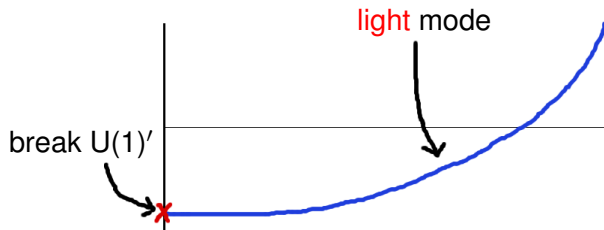
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$$S^{(A')} = \int d^5x \sqrt{-g} \left[-\frac{1}{4} e^{-2\phi} F'_{MN} F'^{MN} \right]$$

- Assume ϕ in coupling has **y-dependent vev**: $\langle \phi \rangle(y) \neq 0$.
- Massless KK mode has **constant profile**, $f^{(0)}(y) = N^{(0)}$:

$$\Rightarrow S^{(A')} \supset ((N^{(0)})^2 \int dy e^{-2\langle \phi \rangle}) \int d^4x \left[-\frac{1}{4} F'^{(0)}_{\mu\nu} F'^{(0)\mu\nu} \right]$$

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$$\hat{A}_M \equiv e^{-\langle \phi \rangle} A'_M \implies S(A') = \int d^5x \left[-\frac{1}{4} \hat{F}_{\mu\nu}^2 - \frac{1}{2} e^{-2ky} \left(\partial_y \hat{A}_\mu \right)^2 - \frac{1}{2} e^{-2ky} (b^2 - 2b) k^2 \hat{A}_\mu^2 - e^{-2ky} b k \hat{A}_\mu^2 (\delta(y) - \delta(y - L)) \right]$$

- \implies 'standard' gauge field in RS but with **bulk and boundary masses**
- \implies **KK decomposition straightforward** (as usual using Bessel functs.)
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Now **break U(1)'** at UV brane by imposing Dirichlet boundary condition.
From boundary conditions, we obtain **mass quantization condition**:

$$\frac{J_b\left(\frac{m_n}{\text{TeV}}\right)}{Y_b\left(\frac{m_n}{\text{TeV}}\right)} = \frac{J_{b-1}\left(\frac{m_n}{k}\right)}{Y_{b-1}\left(\frac{m_n}{k}\right)}$$

Expanding for $m_n \ll \text{TeV}$, we find **exponentially light mode** with mass

$$m_0 \approx \left(\frac{\text{TeV}}{k}\right)^{(b-1)} \text{TeV}$$

⇒ For $k = 10^{18}$ GeV, $b = 1.2$ ($\widehat{f}^{(0)}(y) \propto e^{1.2ky}$) leads to $m_0 \sim \text{GeV}$.

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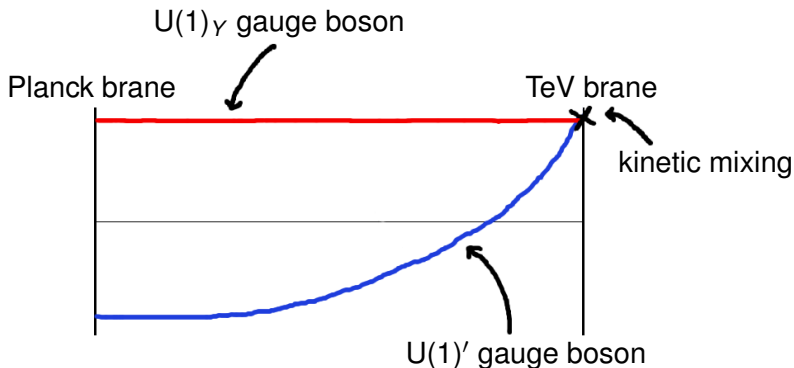
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- Mixing term **mixes KK modes of A_M and A'_M** . Diagonalize action via suitable KK ansatz, mixing term leads to particular **boundary condition**
- For $\langle \phi \rangle \propto y$, KK spectrum etc. straightforward to determine. For weak mixing $\zeta < 1$ find:
 - Spectrum contains **mode with mass $\ll TeV$** (for suitable $\langle \phi \rangle$ and broken $U(1)'$ at UV brane) as well as **massless mode**
 - Light mode and corresponding KK mode tower couple **suppressed by ζ** to $U(1)_Y$ current but unsuppressed to $U(1)'$ current
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DARK MATTER ON THE IR BRANE

- Simplest: Fermion with TeV mass **localized on TeV brane**, charged under $U(1)'$ but neutral under SM gauge group
- Fermion **stable** by virtue of a global $U(1)$
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In 4d: Consider massive Dirac fermion $\chi = (\chi_L, \chi_R)^T$ with small **Majorana mass**. Mass matrix for χ_L and χ_R^\dagger :

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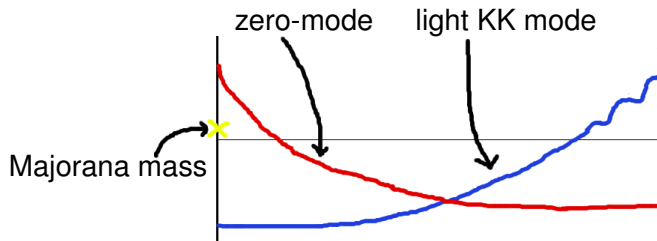
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- Lagrangian for bulk dark matter fermion:

$$\mathcal{L} = - \left[\frac{i}{2} \left(\bar{\chi} \Gamma^M D_M \chi - \bar{D}_M \bar{\chi} \Gamma^M \chi \right) + ic' k \bar{\chi} \chi + i\delta(y) \frac{d'}{2} (\bar{\chi}^c \chi + \text{h.c.}) \right]$$

- Modified boundary condition at UV brane due to Majorana mass. Resulting **mass quantization condition** ($\kappa_{\pm} \equiv c' \pm 1/2$):

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- In warped extra dimension, GeV mass by localizing $U(1)'$ gauge boson away from Planck brane and then breaking $U(1)'$ at Planck brane
- Gauge boson can be localized if kinetic term has form $e^{-2\langle\phi\rangle} F_{MN}^2$ with y -dependent vev $\langle\phi\rangle$
- Case $\langle\phi\rangle \propto y$ easy to analyze. Checked also presence of light mode for case $\langle\phi\rangle \propto e^{-ay}$. Showed how to obtain such vev.
- Small mass split for dark matter for bulk fermion with Majorana mass term on UV brane (useful to avoid direct detection constraints and to reconcile DAMA with other experiments via inelastic dark matter scenario)

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